

# BLACK HOLE CHEMISTRY

Robert Mann

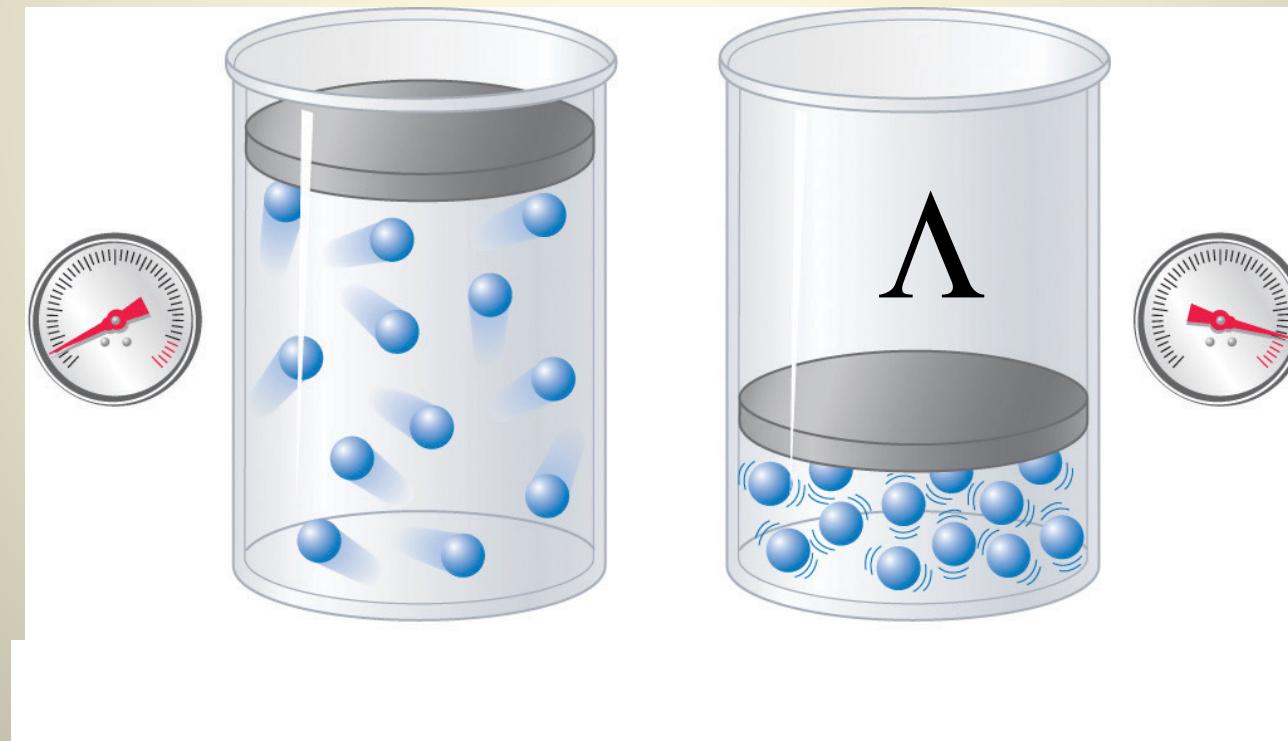


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<https://boardgamegeek.com/image/2241156/alchemists>

# Black Hole Chemistry: Pressure from the Vacuum

$$p = -\frac{\Lambda}{8\pi G}$$



# Recent Results in Black Hole Chemistry

- Accelerating Black Hole Thermodynamics
- Holographic Black Hole Chemistry
  - Holographic Smarr Relation
  - Holographic Complexity and Thermodynamic Volume
  - Central Charge Criticality
- Black Hole Multicriticality
- And more!

# Accelerating Black Holes

- Basic features
  - both an event horizon and an acceleration horizon
  - Thermodynamic equilibrium can't be maintained
- Asymptotically AdS
  - Acceleration horizon removed for small acceleration
  - Cosmic string suspends BH away from the centre of AdS
- Conflicting Results for Thermodynamics
  - Differing identifications of mass Appels/Gregory/ Kubiznak PRL
  - Role of conical deficits in first law not clear 117 (2016) 131303;  
JHEP 1705 (2017)116
  - Free energy/action not compatible Astorino PRD95 (2017) 064007  
Anabalon/Appels/Gregory/  
Kubiznak/Mann/Ovgun PRD

# Metric for ABH in AdS

$$ds^2 = \frac{1}{\Omega^2} \left[ -fdt^2 + \frac{dr^2}{f} + r^2 \left( \frac{d\theta^2}{\Sigma} + \Sigma \sin^2 \theta \frac{d\phi^2}{K^2} \right) \right]$$

$$\Omega = 1 + Ar \cos \theta$$

$$\Sigma = 1 + 2mA \cos \theta$$

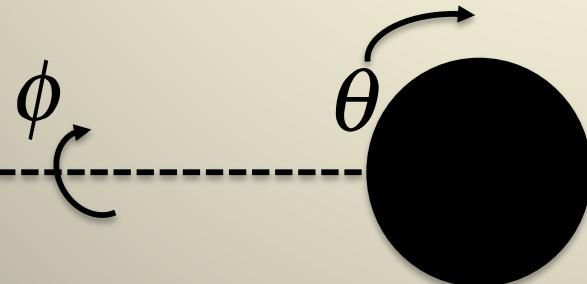
Hong/Teo

$$f(r) = (1 - A^2 r^2) \left( 1 - \frac{2m}{r} \right) + \frac{r^2}{\ell^2}$$

CQG **20** (2003) 3269

String  
Tension

$$\mu_{\pm} = \frac{\delta_{\pm}}{8\pi} = \frac{1}{4} \left( 1 - \frac{\Sigma(\theta_{\pm})}{K} \right) = \frac{1}{4} \left( 1 - \frac{1 \pm 2mA}{K} \right)$$



Parameterizes  
deficit angle

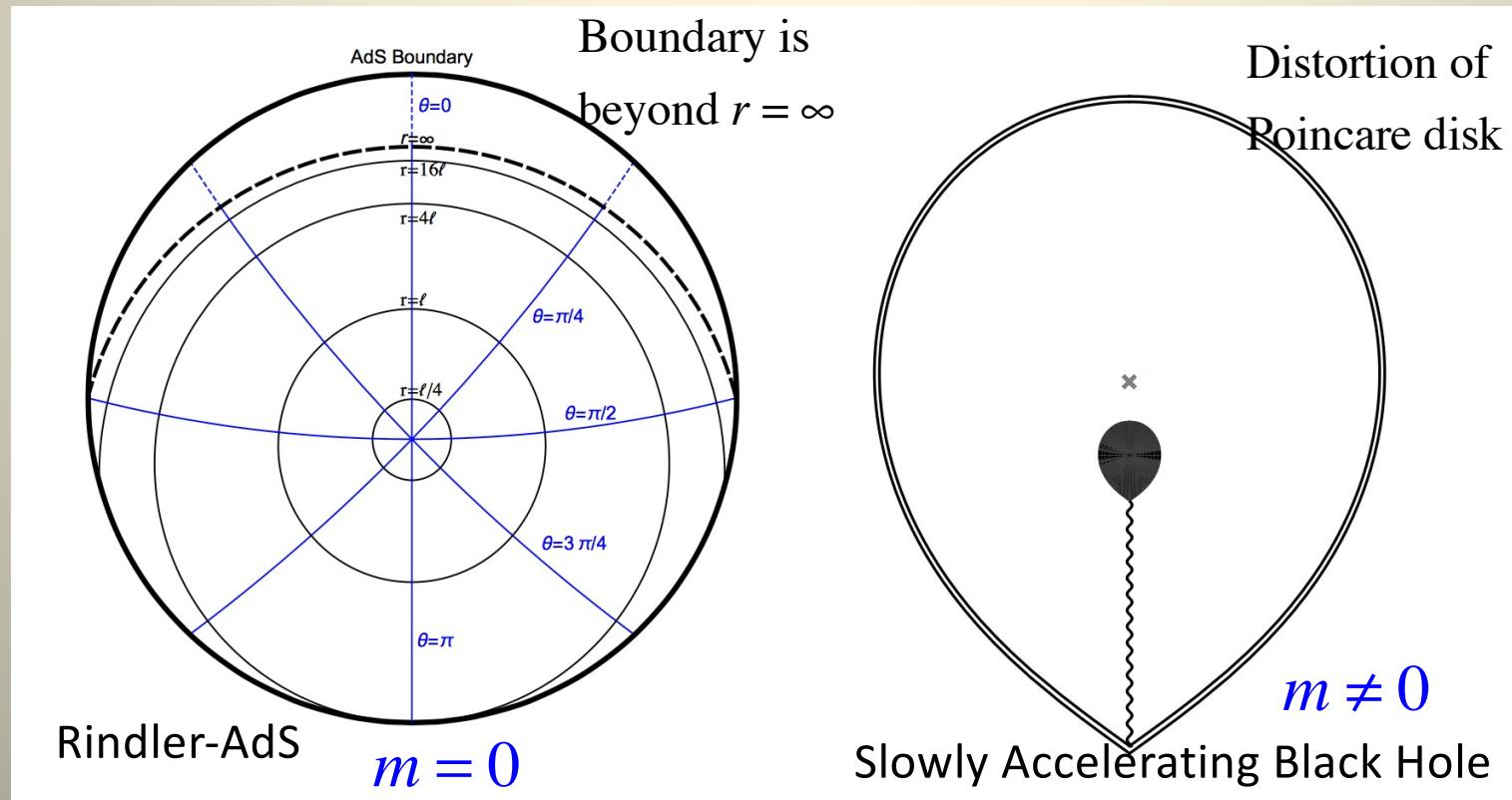
# Metric for ABH in AdS

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$$f(r) = (1 - A^2 r^2) \left( 1 - \frac{2m}{r} \right) + \frac{r^2}{\ell^2}$$

$$\Omega = 1 + Ar \cos \theta$$

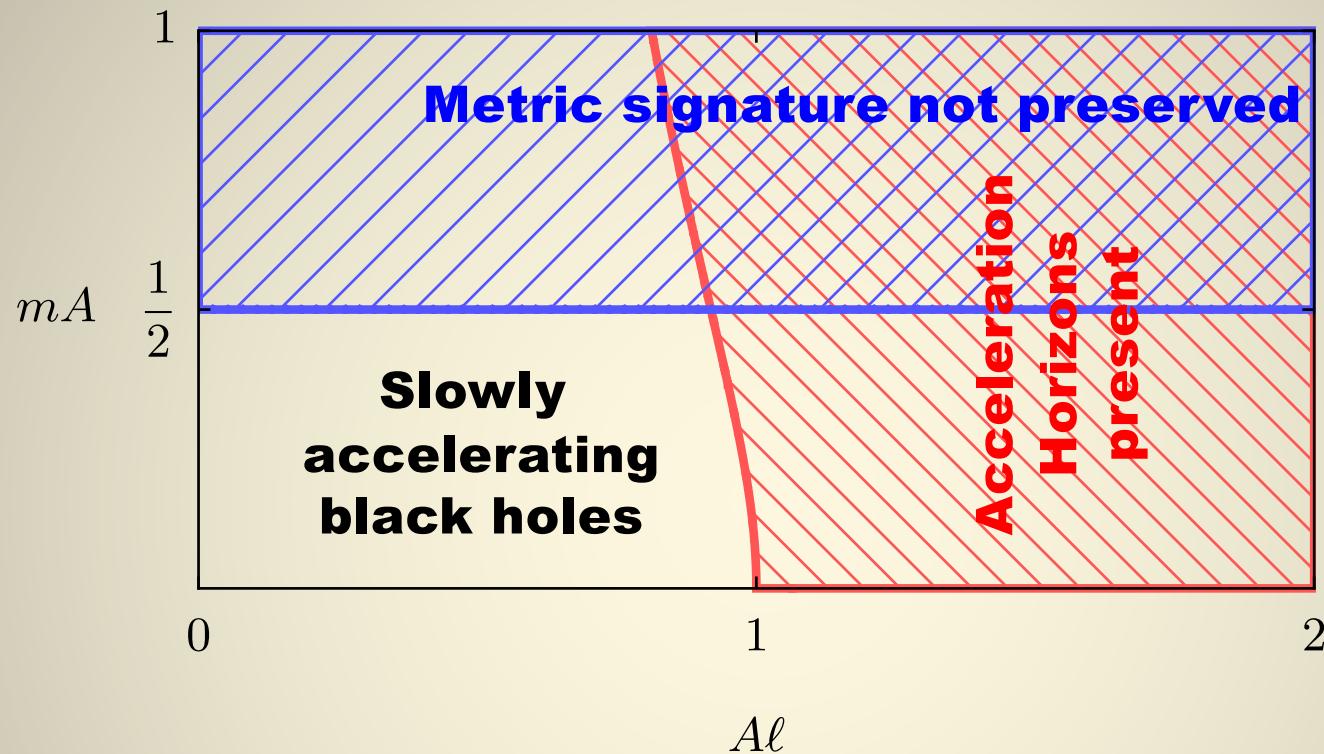
$$\Sigma = 1 + 2mA \cos \theta$$



$$A\ell < \frac{3\sqrt{6}}{8}$$

No Acceleration Horizons

$$f(-1/A \cos \theta) > 1$$



$$ds^2 = \frac{1}{\Omega^2} \left[ -fdt^2 + \frac{dr^2}{f} + r^2 \left( \frac{d\theta^2}{\Sigma} + \Sigma \sin^2 \theta \frac{d\phi^2}{K^2} \right) \right]$$

$$f(r) = (1 - A^2 r^2) \left( 1 - \frac{2m}{r} \right) + \frac{r^2}{\ell^2}$$

$$\begin{aligned}\Omega &= 1 + Ar \cos \theta \\ \Sigma &= 1 + 2mA \cos \theta\end{aligned}$$

# Asymptotics

$$\Omega = 1 + A r \cos \theta$$

$$\Sigma = 1 + 2m A \cos \theta$$

$$ds^2 = \frac{1}{\Omega^2} \left[ -fdt^2 + \frac{dr^2}{f} + r^2 \left( \frac{d\theta^2}{\Sigma} + \Sigma \sin^2 \theta \frac{d\phi^2}{K^2} \right) \right]$$

$$1 + \frac{R^2}{\ell^2} = \frac{1 + (1 - A^2 \ell^2)r^2 / \ell^2}{(1 - A^2 \ell^2)\Omega^2}$$

$$R \sin \vartheta = \frac{r \sin \theta}{\Omega}$$

$m = 0$

$$ds_{AdS}^2 = - \left( 1 + \frac{R^2}{\ell^2} \right) \alpha^2 dt^2 + \frac{dR^2}{1 + \frac{R^2}{\ell^2}} + R^2 \left( d\vartheta^2 + \sin^2 \vartheta \frac{d\phi^2}{K^2} \right)$$

$$\alpha = \sqrt{1 - A^2 \ell^2}$$

Correct time coordinate  
is  $\tau = \alpha t$

# Thermodynamic Quantities

Temperature

$$T = \frac{f'(r_+)}{4\pi\alpha} = \frac{1 + 3\frac{r_+^2}{\ell^2} - A^2 r_+^2 \left( 2 + \frac{r_+^2}{\ell^2} - A^2 r_+^2 \right)}{4\pi\alpha r_+ (1 - A^2 r_+^2)}$$

Entropy

$$S = \frac{\mathcal{A}}{4} = \frac{\pi r_+^2}{K(1 - A^2 r_+^2)}$$

Mass

- Thermodynamic
  - Via the Smarr Relation
- Conformal
  - Via Electric part of Weyl Tensor
- Holographic
  - Via AdS/CFT counterterms

$$f(r_+) = (1 - A^2 r_+^2) \left( 1 - \frac{2m}{r_+} \right) + \frac{r_+^2}{\ell^2} = 0$$

$$M = \frac{\alpha m}{K} = \sqrt{1 - A^2 \ell^2} \frac{m}{K}$$

# Action

$$I = \frac{\beta}{2\alpha K} \left( m - 2mA^2\ell^2 - \frac{r_+^3}{\ell^2(1-A^2r_+^2)^2} \right)$$

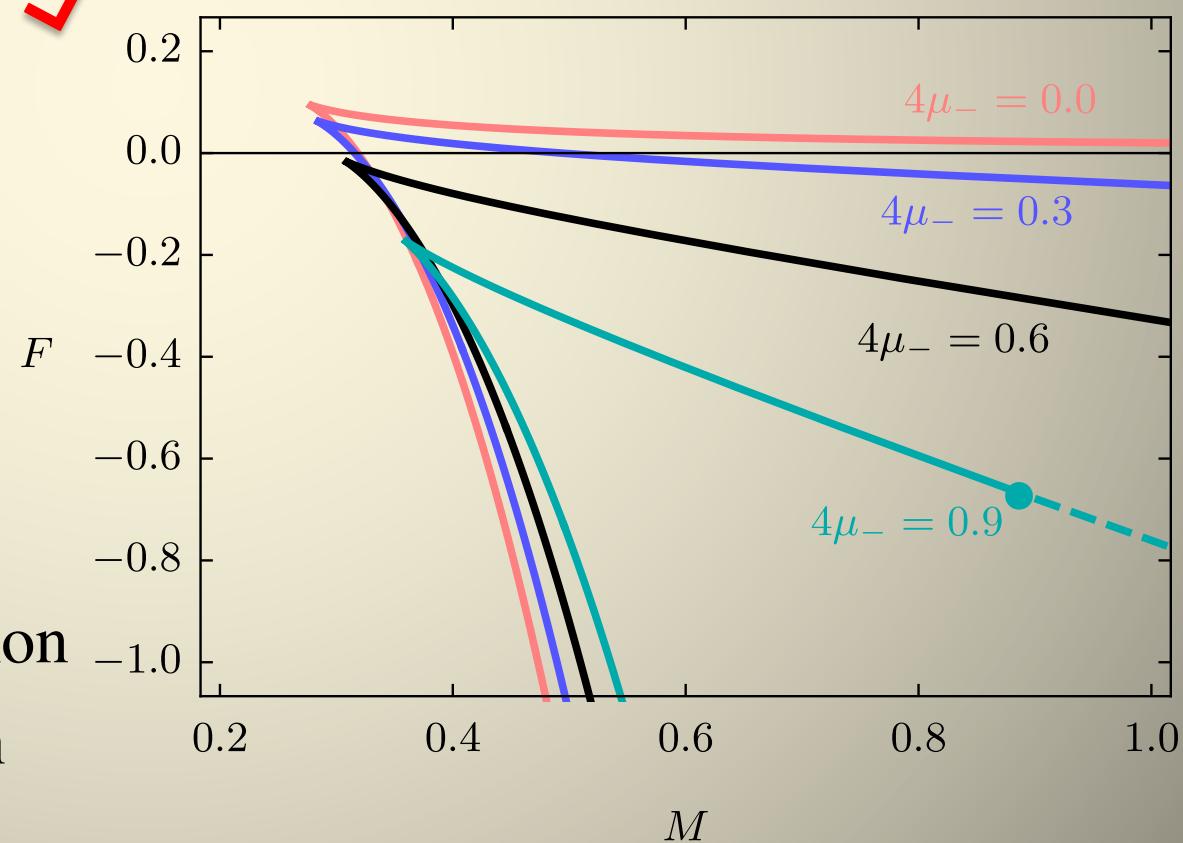
$$\beta = 1/T$$

$\rightarrow F = I / \beta = M - TS \checkmark$

$$\mu_{\pm} = \frac{\delta_{\pm}}{8\pi} = \frac{1}{4} \left( 1 - \frac{1 \pm 2mA}{K} \right)$$

$\mu_- = 0$ : Hawking-Page transition

$\mu_- > 0$ : No clear interpretation



# With Charge and Rotation

$$ds^2 = \frac{1}{H^2} \left\{ -\frac{f(r)}{\Sigma} \left[ \frac{dt}{\alpha} - a \sin^2 \theta \frac{d\varphi}{K} \right]^2 + \frac{\Sigma}{f(r)} dr^2 + \frac{\Sigma r^2}{h(\theta)} d\theta^2 + \frac{h(\theta) \sin^2 \theta}{\Sigma r^2} \left[ \frac{adt}{\alpha} - (r^2 + a^2) \frac{d\varphi}{K} \right]^2 \right\}$$

$$F = dB \quad B = -\frac{e}{\Sigma r} \left[ \frac{dt}{\alpha} - a \sin^2 \theta \frac{d\varphi}{K} \right] + \Phi_t dt$$

$$M = \frac{m(\Xi + a^2 / \ell^2)(1 - A^2 \ell^2 \Xi)}{K \Xi \alpha (1 + a^2 A^2)} \quad T = \frac{f'_+ r_+^2}{4\pi \alpha (r_+^2 + a^2)}$$

$$S = \frac{\pi(r_+^2 + a^2)}{K(1 - A^2 r_+^2)} \quad Q = \frac{e}{K} \quad \Phi = \Phi_t = \frac{er_+}{(r_+^2 + a^2)\alpha}$$

$$J = \frac{ma}{K^2} \quad \Omega = \Omega_H - \Omega_\infty = \left( \frac{Ka}{\alpha(r_+^2 + a^2)} \right) - \left( -\frac{aK(1 - A^2 \ell^2 \Xi)}{\ell^2 \Xi \alpha (1 + a^2 A^2)} \right)$$

$$P = \frac{3}{8\pi \ell^2} \quad V = \frac{4\pi}{3K\alpha} \left[ \frac{r_+ (r_+^2 + a^2)}{(1 - A^2 r_+^2)^2} + \frac{m[a^2(1 - A^2 \ell^2 \Xi) + A^2 \ell^4 \Xi (\Xi + a^2 / \ell^2)]}{(1 + a^2 A^2) \Xi} \right]$$

Anabalon/Appels/  
Gray/Gregory/

Kubiznak/Mann/Ovgun  
JHEP 1904 (2019) 096

$$\Sigma = 1 + \frac{a^2}{r^2} \cos^2 \theta$$

$$H = 1 + Ar \cos \theta$$

$$h(\theta) = 1 + 2mA \cos \theta$$

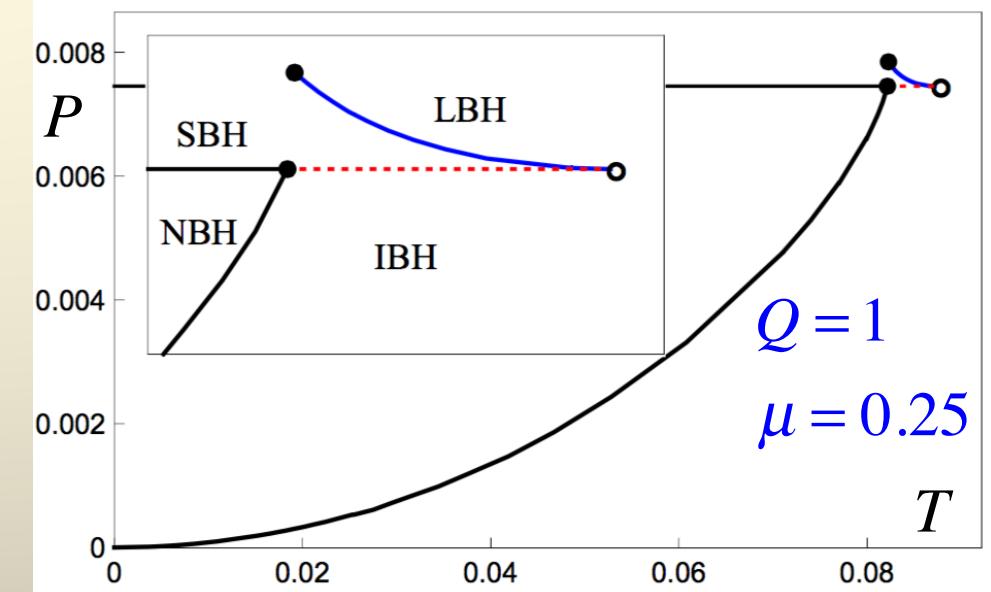
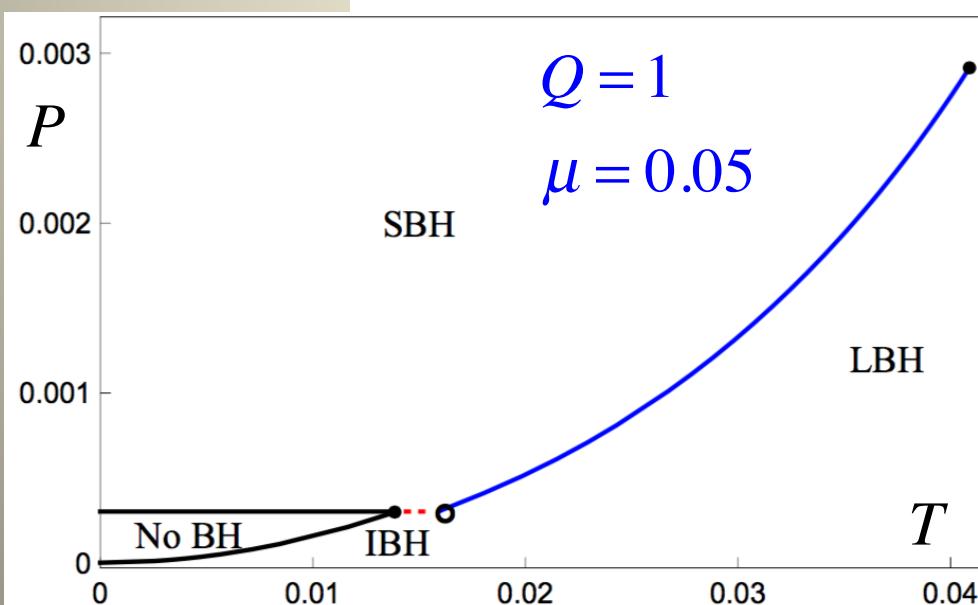
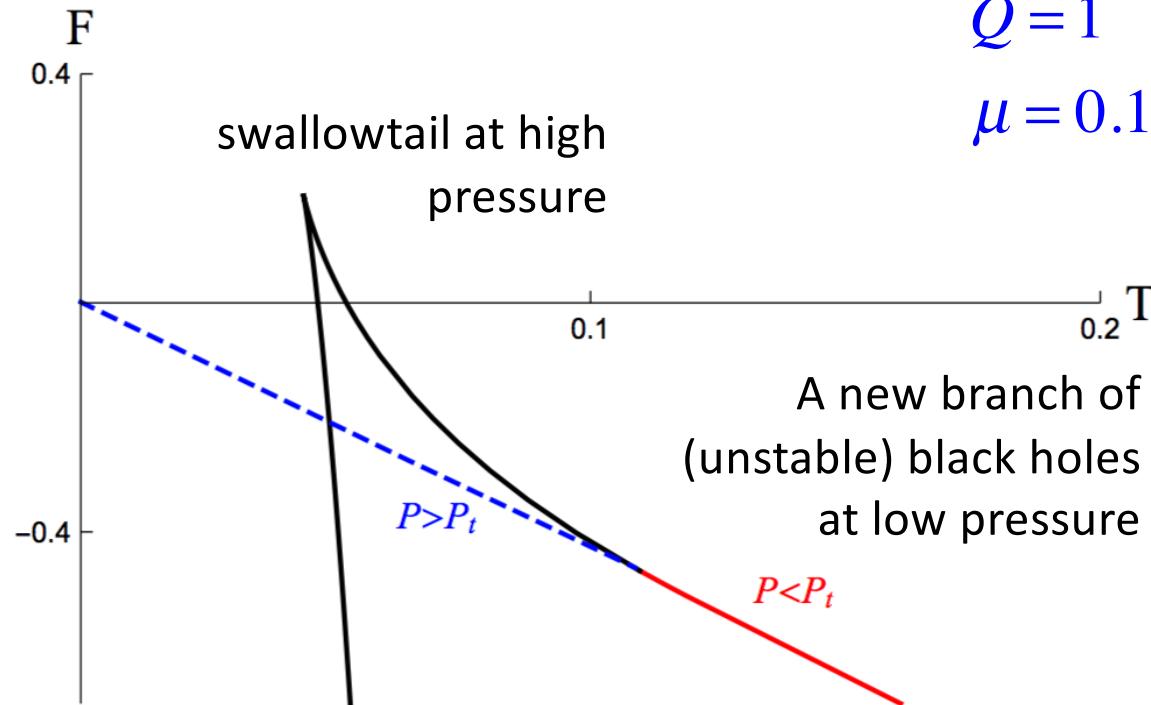
$$+ \left[ A^2(a^2 + e^2) - \frac{a^2}{\ell^2} \right] \cos^2 \theta$$

$$f(r) = (1 - A^2 r^2) \left[ 1 - \frac{2m}{r} + \frac{a^2 + e^2}{r^2} \right] + \frac{r^2 + a^2}{\ell^2}$$

$$\Xi = 1 - \frac{a^2}{\ell^2} + A^2(e^2 + a^2)$$

First law and  
Smarr OK!

# Snapping Swallowtails



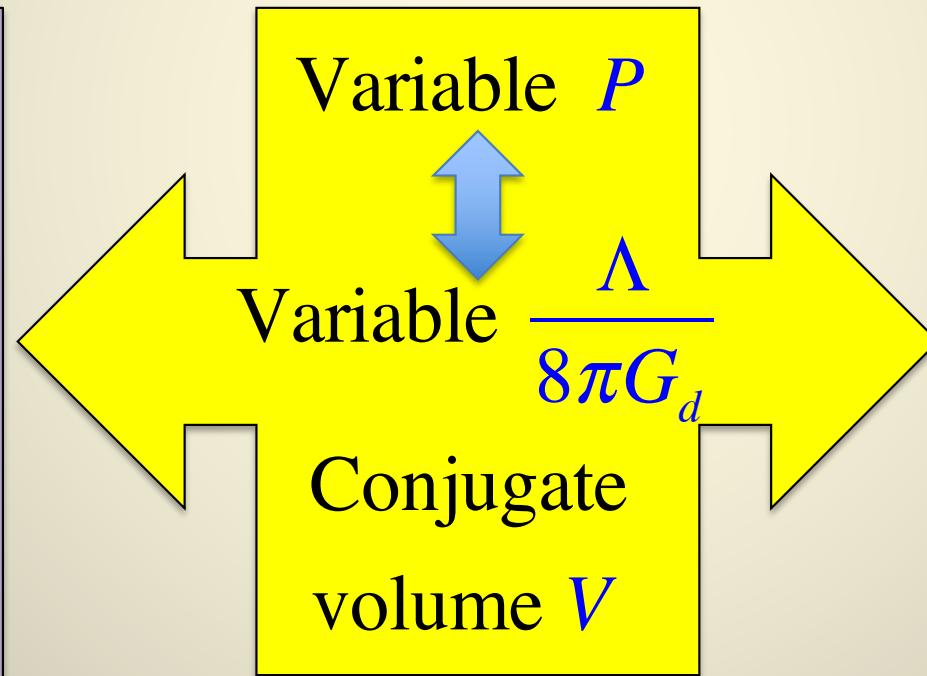
# Variable $\Lambda$ and AdS/CFT?

$$AdS_5 \times S^5 \Leftrightarrow \mathcal{N} = 4 \text{ SU(N) Super Yang-Mills}$$

near-horizon geometry of  $N$  coincident D3  
branes in type IIB supergravity

$$l^4 = \frac{\sqrt{2}\ell_{10-Pl}^4}{\pi^2} N$$

Vary  $N^2$   
  
Conjugate  
chemical  
potential  
  
for colour  $\mu$



Variable Field  
Theory Volume  
  
Conjugate  
Gravitational  
Coupling  $G_d$

Johnson CQG **31** (2014) 205002

Dolan JHEP **1014** (2014) 179

Kastor/Ray/Traschen JHEP **1114** (2014) 120

Karch/Robinson JHEP **1015** (2015) 73

# Holographic Smarr Relation

Karch/Robinson

JHEP 1512 (2015) 073

$$\ell^4 = \frac{\sqrt{2} \ell_{Pl}^4}{\pi^2} N \quad \xrightarrow{\text{large } N} \quad \Omega(N, \mu, T, l) = N^2 \Omega_0(\mu, T, l)$$

Free  
energy

$$E = TS + \tilde{\phi} \tilde{Q} + \Omega J + \tilde{\mu} C$$

$$E = (D - 2)pV \quad \text{CFT Eqn of State}$$

polynomial functions

$$\Omega(N, \mu, T, \alpha_j, R) = \sum_{k=0} g_k(N) \Omega^k(\mu, T, \alpha_j, R)$$

Sinamuli/Mann

PRD 96 (2017) 086008

Extension to  
Lovelock Gravity

$$\text{scaling: } \alpha_k = [L]^{D-2} \quad \xrightarrow{\hspace{10em}}$$

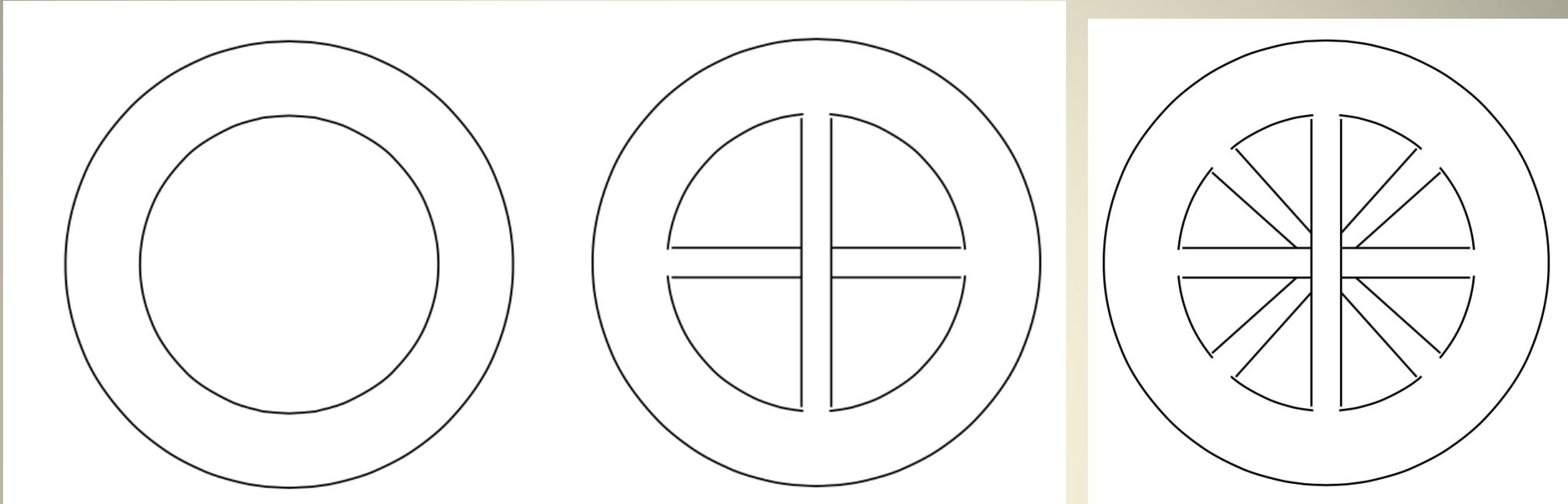
$$g_k(N) = \beta_k (\alpha_k)^{\frac{D-2}{2(k-1)}}$$

$$\ell \frac{\partial}{\partial \ell} + \sum_{k=1} 2(k-1) \alpha_k \frac{\partial}{\partial \alpha_k} = (d-2) \sum_{k=0} g_k \frac{\partial}{\partial g_k} \quad \Omega = \sum_{k=0} g_k \frac{\partial \Omega}{\partial g_k}$$

$$\frac{D-3}{D-2} M = TS - \frac{2}{D-2} PV + \sum_{k=2}^K \frac{2(k-1)}{D-2} \Psi^k \alpha_k + \frac{D-3}{D-2} \sum_j \Phi^j Q^j$$

# Evaluating $g_k(N)$

Sinamuli/Mann  
PRD **96** (2017) 086008



$$N^2(g_{YM})^0 = N^2 \lambda^0$$

$$\rightarrow \mathcal{L}^{(0)}$$

$$N^2(g_{YM})^4 = N^0 \lambda^2$$

$$\rightarrow \mathcal{L}^{(1)}$$

$$N^2(g_{YM})^8 = N^{-2} \lambda^4$$

$$\rightarrow \mathcal{L}^{(2)}$$

$$L^{(k)} \sim R^k \rightarrow N^{2(1-k)}$$

# Holographic Complexity

$\mathcal{C}(\Sigma)$ : quantum complexity of boundary state at time slice  $\Sigma$

**Complexity = Volume**

$$\mathcal{C}_V(\Sigma) = \max_{\Sigma=\partial\mathcal{B}} \left[ \frac{\mathcal{V}(\mathcal{B})}{G_N R} \right]$$

$\mathcal{V}(\mathcal{B})$ : volume of extremal co-dim-1 slice  $\mathcal{B}$

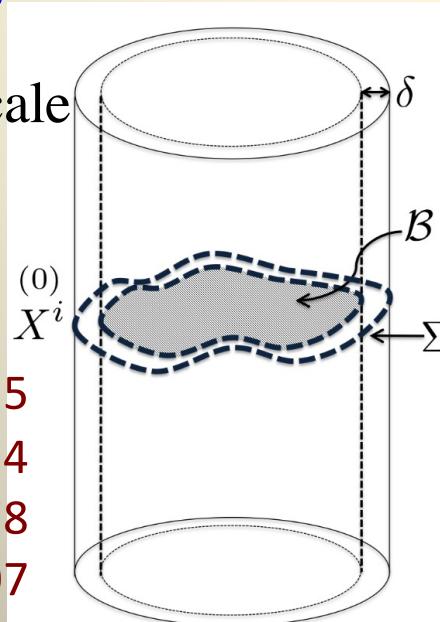
$R$ : arbitrary length scale (eg AdS length)

Susskind 1403.5695

Fortsch. Phys. 64 (2016) 24

Stanford/Susskind 1406.2678

PRD 90 (2014) 126007



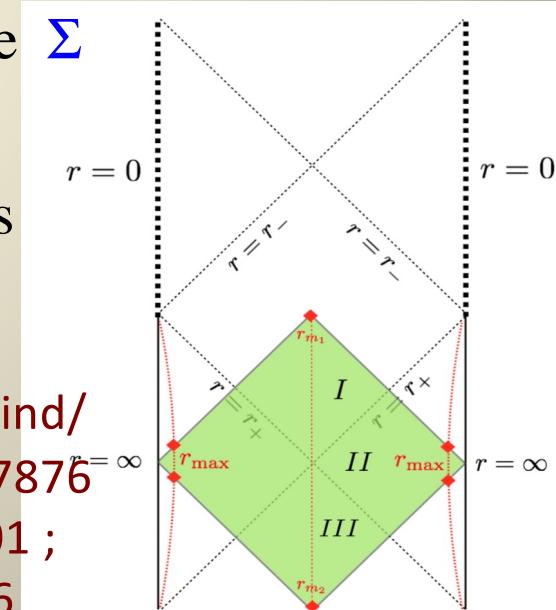
**Complexity= Action**

$$\mathcal{C}_A(\Sigma) = \frac{I_{WDW}}{\pi}$$

$WDW$  : Wheeler-de Witt patch  
domain of dependence of  
Cauchy slice  $\Sigma$

$I_{WDW}$  : action on this domain

Brown/Roberts/Susskind/  
Swingle/Zhao 1509.07876  
PRL 116 (2016) 191301 ;  
PRD 93 (2016) 086006



# Complexity and Black Hole Thermodynamics

## Complexity of Formation

Lehner/Myers/Poisson/Sorkin  
PRD **94** (2016) 084046

Complexity = Volume

$$\Delta \mathcal{C}_V(\Sigma) = \lim_{r_{\max} \rightarrow \infty} \frac{[\mathcal{V}(\Sigma) - 2\mathcal{V}_{AdS}]}{G_N R}$$

Complexity= Action

$$\Delta \mathcal{C}_A(\Sigma) = \frac{I_{WDW} - 2I_{AdS}}{\pi}$$

General Expectation (near extremality)

$$\Delta \mathcal{C}_V \sim \tilde{k}_d S \log \frac{\mu}{T}$$

$$\Delta \mathcal{C}_A \sim k_d S \log \frac{\mu}{T}$$

Carmi/Chapman/Marrochio  
Myers/Sugishita  
JHEP **2017** (2017)

Chapman/Heller/Marrochio/Pastawski  
PRL **120** (2017) 121602

Based on calculations for charged spherically symmetric blackholes

# Complexity and Black Hole Thermodynamics

## Complexity of Formation

Complexity = Volume

$$\Delta\mathcal{C}_V(\Sigma) = \lim_{r_{\max} \rightarrow \infty} \frac{[\mathcal{V}(\Sigma) - 2\mathcal{V}_{AdS}]}{G_N R}$$

We find

$$\Delta\mathcal{C}_V \sim S \log \frac{\mu}{T}$$

$$\Delta\mathcal{C}_V \sim V^{\frac{D-2}{D-1}} \log \frac{\Omega_H}{T}$$

Complexity= Action

$$\Delta\mathcal{C}_A(\Sigma) = \frac{I_{WDW} - 2I_{AdS}}{\pi}$$

Balushi /Hennigar/Khunduri  
PRL 126 (2021) 101601  
JHEP 2105 (2021) 226

$$\Delta\mathcal{C}_A \sim S \log \frac{\mu}{T}$$

$$\Delta\mathcal{C}_A \sim V^{\frac{D-2}{D-1}} \log \frac{\Omega_H}{T}$$

where  $V$  is the thermodynamic volume of a rotating black hole

# Complexity $\sim$ Thermodynamic Volume

## 3+1 Dimensional Charged Black Holes

$$\text{Entropy } S = \pi r_+^2$$

$$\text{Volume } V = \frac{4\pi r_+^3}{3}$$

## (Most) D-Dimensional Spherical Black Holes

$$\text{Entropy } S \sim r_+^{D-2}$$

$$\text{Volume } V \sim r_+^{D-1}$$

- Degeneracy between entropy and volume for most spherically symmetric black holes
- This degeneracy is broken for rotating black holes

$$\text{Entropy } S \sim \pi(r_+^2 + a^2)$$

$$\text{Volume } V = \frac{4\pi l^2 r_+}{3} \frac{a^2 + r_+^2}{l^2 - a^2} + 4\pi a J / 3$$

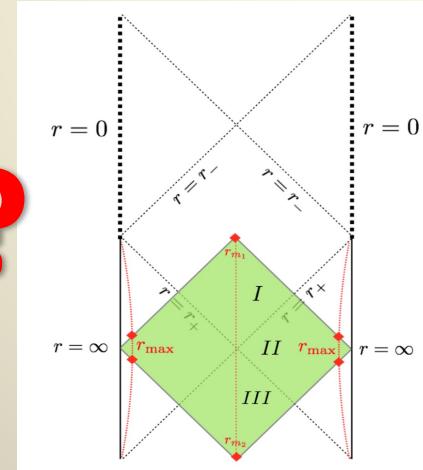
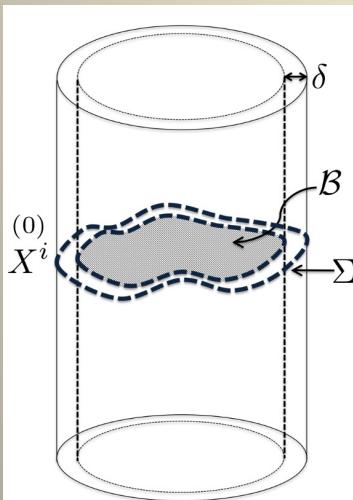
Cvetic/Gibbons/Kubiznak/Pope

PRD84 (2011) 024037

# Computing Complexity in Rotating Black Holes?

- Null hypersurfaces are complicated functions of polar angle as well as radial variable
- Not clear how to match null hypersurfaces from AdS boundary to interior

AlBalushi/Mann  
CQG 36 (2019) 245017



Problem drastically simplifies for multiply-rotating black holes in odd-dimensions with equal angular momenta

- Caustics can form if charge is sufficiently large

Imseis/AlBalushi/Mann  
CQG 38 (2021) 045018

# Multiply Rotating AdS Black Holes in D=2N+3 Dimensions

$$ds^2 = -f(r)^2 dt^2 + g(r)^2 dr^2 + r^2 \hat{g}_{ab} dx^a dx^b$$

$$+ h(r)^2 \left[ d\psi + A(x) - \Omega(r) dt \right]^2$$

$$g(r)^2 = \left( 1 + \frac{r^2}{\ell^2} - \frac{2m\Xi}{r^{2N}} + \frac{2ma^2}{r^{2N+2}} \right)^{-1} \quad f(r) = \frac{r}{g(r)h(r)}$$

$$h(r)^2 = r^2 \left( 1 + \frac{2ma^2}{r^{2N+2}} \right) \quad \Omega(r) = \frac{2ma}{r^{2N} h^2} \quad \Xi = 1 - \frac{a^2}{\ell^2}$$

Gibbons/Lu/Page/Pope [hep-th/0404008]

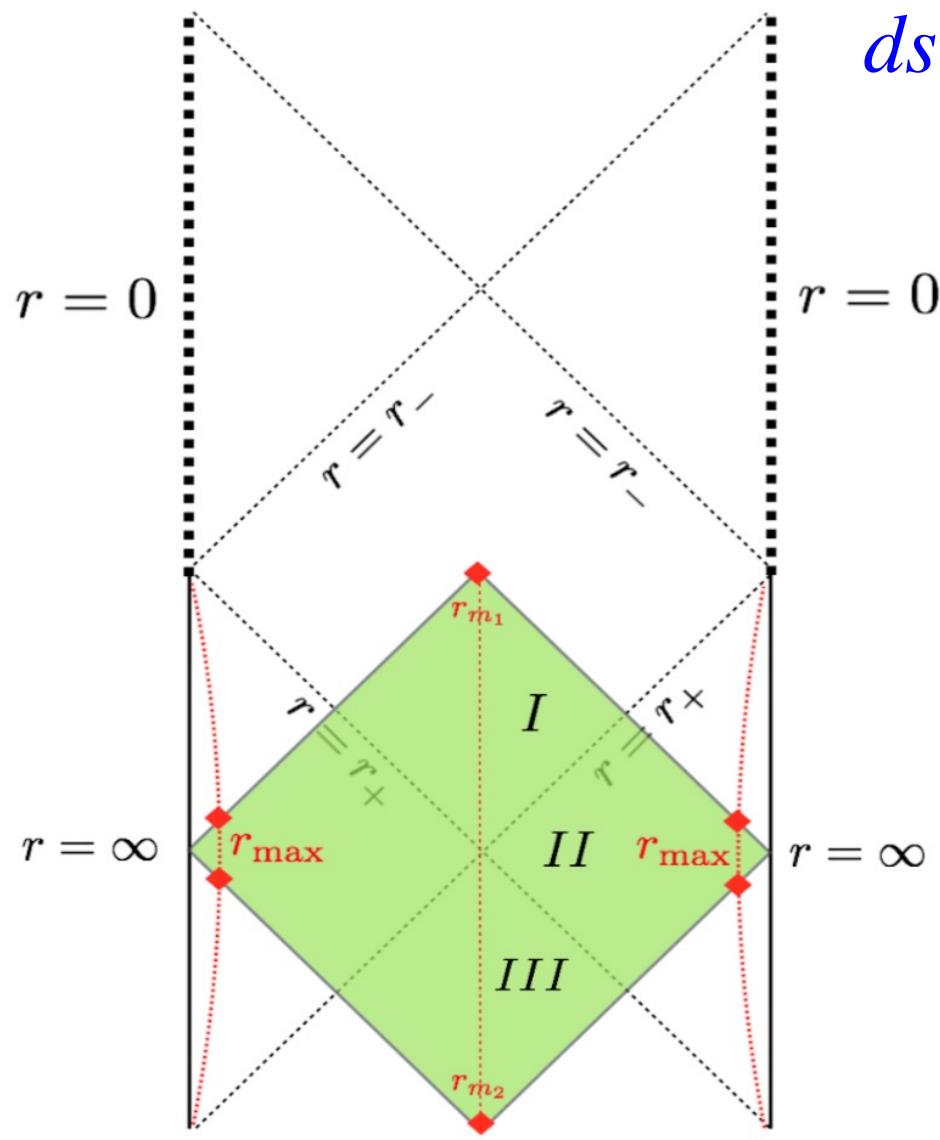
J Geom Phys (2005) 49

Kunduri/Lucietti /Reall

[hep-th/0606076] PRD74 (2006) 084021

$$\hat{R}_{ij} = 2(N+1)\hat{g}_{ij} \quad \left\{ \begin{array}{l} D=5 \\ \hat{g} = \frac{1}{4}(d\theta^2 + \sin^2 \theta d\phi^2) \\ A = \frac{1}{2}\cos \theta d\phi \end{array} \right.$$

# Causal Structure



$$ds^2 = -f(r)^2 dt^2 + g(r)^2 dr^2 + r^2 \hat{g}_{ab} dx^a dx^b + h(r)^2 [d\psi + A - \Omega(r)dt]^2$$

Simple null hypersurfaces!

Horizon Generator

$$\xi = \frac{\partial}{\partial t} + \Omega_H \frac{\partial}{\partial \psi}$$

$$\Omega_H = \frac{2ma}{r_+^{2N+2} + 2ma^2}$$

# Complexity equals Action

$$\pi \Delta \mathcal{C}_A(Y) = I_{\text{WDW}} - 2I_{\text{AdS}}$$

$$= \frac{\Lambda \Omega_{2N+1}}{2(N+1)(2N+1)\pi G_N} \left[ \int_{r_{m_0}}^{\infty} dr \ r^{2N+1} g(r)^2 h(r) \right.$$

$$\left. - \int_0^{\infty} dr \frac{r^{2(N+1)}}{1+r^2/\ell^2} \right]$$

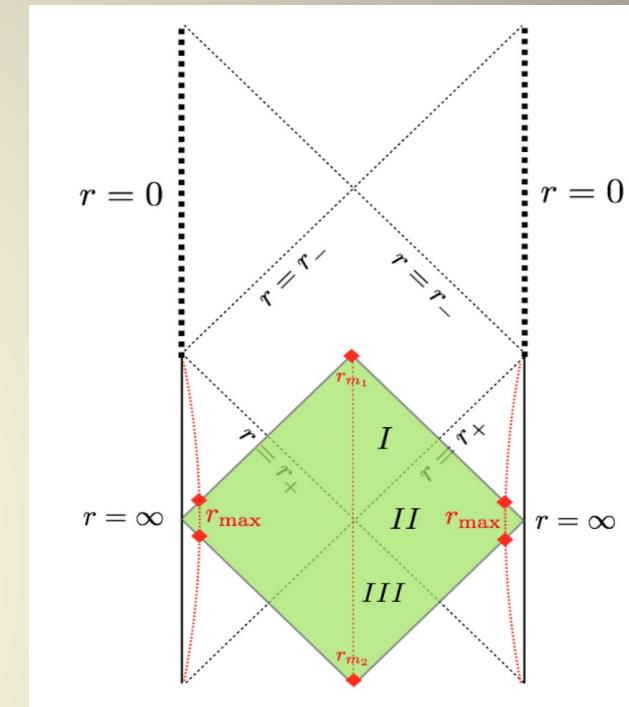
$$- \frac{\Omega_{2N+1}(r_{m_0})^{2N+1}}{2\pi G_N(2N+1)} - \frac{\Omega_{2N+1}}{2\pi G_N} \int_{r_{m_0}}^{\infty} dr \ r^{2N} \left[ h(r) \frac{\Theta'}{\Theta} + 1 \right]$$

$$- \frac{\Omega_{2N+1}}{4\pi G_N} (r_{m_0})^{2N} h(r_{m_0}) \log \ell_{\text{ct}}^2 \Theta(r_{m_0})^2 |f(r_{m_0})^2|$$

Balushi /Hennigar/Khunduri  
 PRL 126 (2021) 101601  
 JHEP 2105 (2021) 226

We find

$$\mathcal{C}_A(\Sigma) = \frac{I_{\text{WDW}}}{\pi}$$

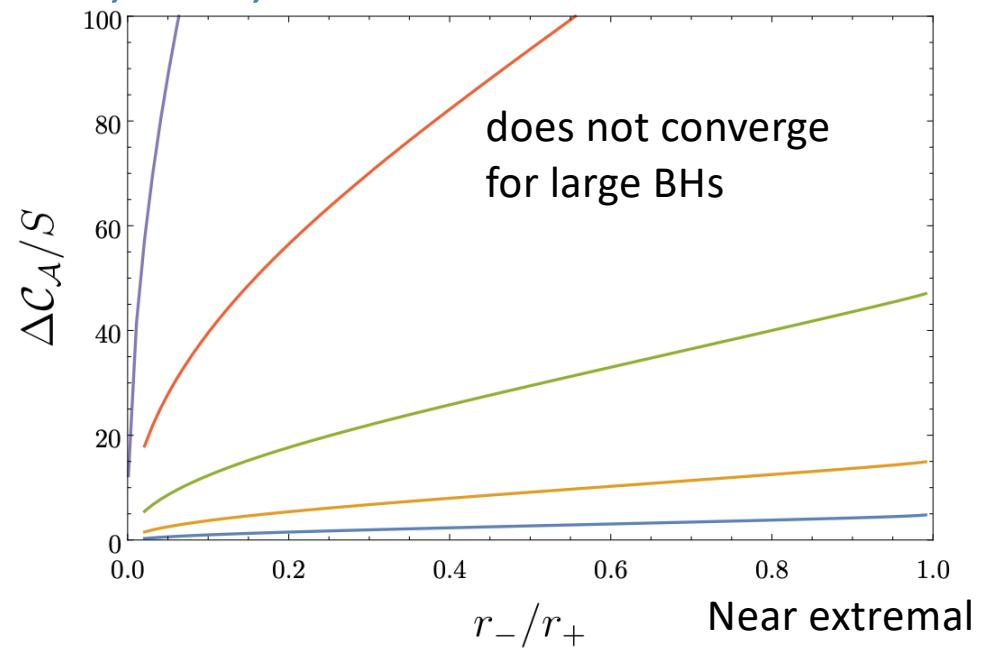
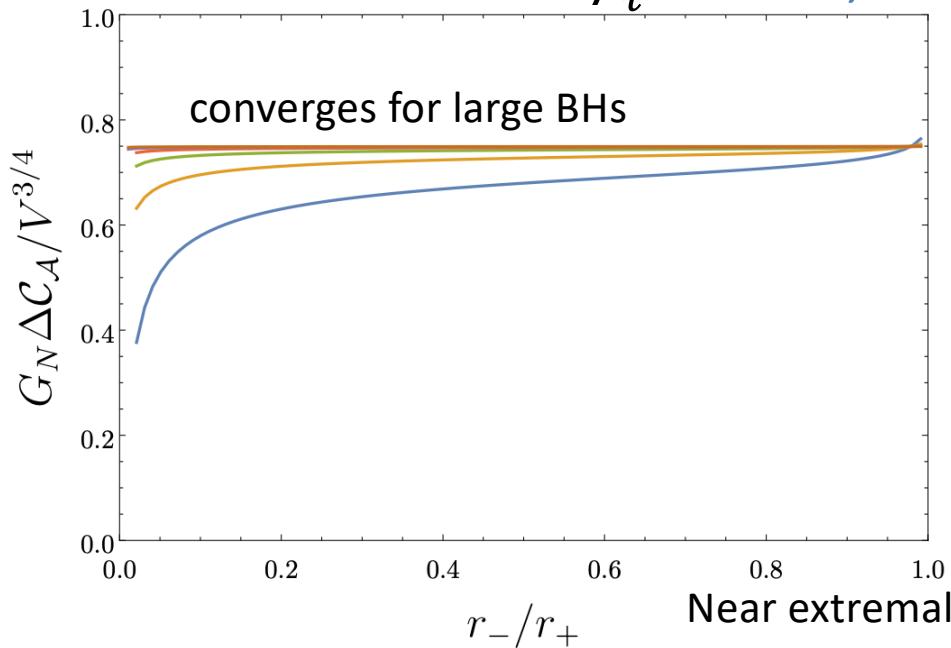


$$\Theta = \frac{1}{f(r)g(r)} \left[ \frac{2N}{r} + \frac{h'}{h} \right]$$

$$\Delta \mathcal{C}_A(\Sigma) = \Sigma_g C_T \left( \frac{V}{V_{\text{AdS}}} \right)^{\frac{D-2}{D-1}}$$

D=5

$$r_+/\ell = 10^3, 10^4, 10^5, 10^6, 10^7$$

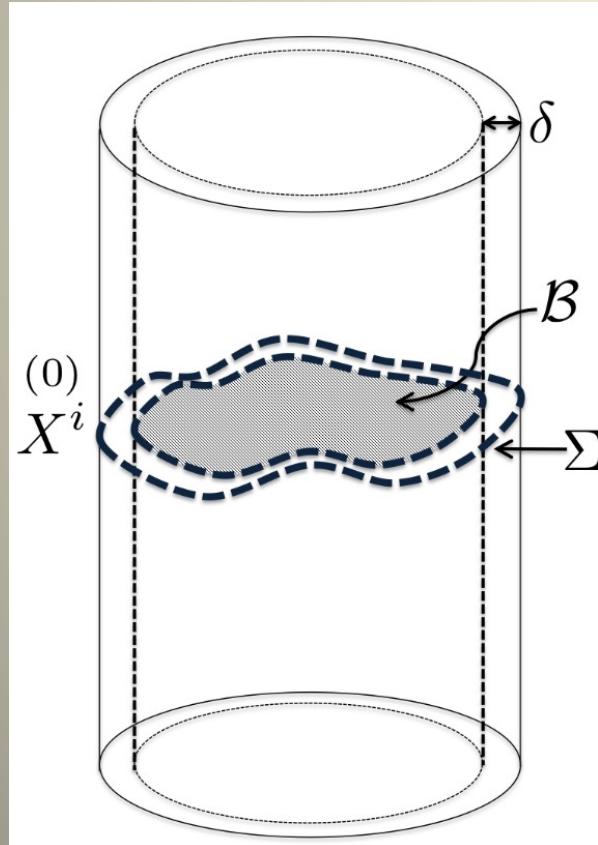


$$\Delta \mathcal{C}_v = \sum_g C_T \left( \frac{V}{V_{AdS}} \right)^{\frac{D-2}{D-1}} = \sum_g C_T \left( \frac{V}{V_{AdS}} \right)^{\frac{3}{4}}$$

$$\neq KS$$

# Complexity equals Volume

$$\mathcal{C}_V(\Sigma) = \max_{\Sigma = \partial \mathcal{B}} \left[ \frac{\mathcal{V}(\mathcal{B})}{G_N R} \right]$$



$$\mathcal{V} = [\mathcal{V} - 2\mathcal{V}_{AdS}]$$

$$= 2\Omega_{D-2} \int_{r_+}^{r_{\max}} dr r^{(D-3)} h(r) g(r)$$

$$- 2\Omega_{D-2} \int_0^{r_{\max}} \frac{r^{D-2} dr}{\sqrt{1 + r^2 / \ell^2}}$$

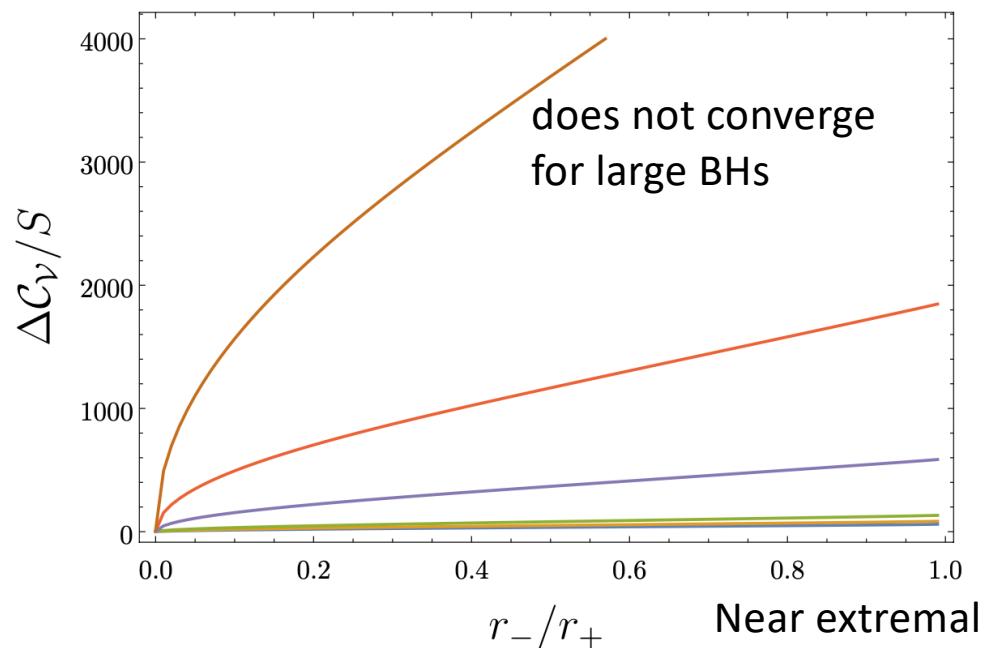
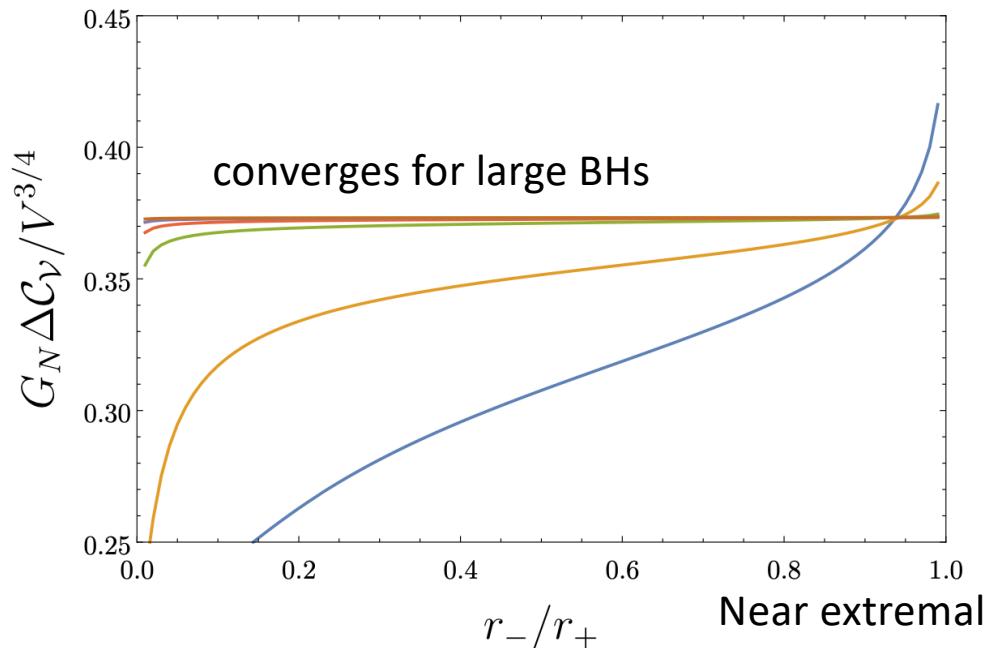
Balushi /Hennigar/Khunduri  
 PRL **126** (2021) 101601  
 JHEP **2105** (2021) 226

We find

$$\Delta \mathcal{C}_V = \lim_{r_{\max} \rightarrow \infty} \frac{[\mathcal{V} - 2\mathcal{V}_{AdS}]}{G_N R} = \Sigma_g C_T \left( \frac{V}{V_{AdS}} \right)^{\frac{D-2}{D-1}}$$

D=5

$$r_+/\ell = 10, 10^2, 10^3, 10^4, 10^5, 10^6$$



$$\Delta C_V = \sum_g C_T \left( \frac{V}{V_{\text{AdS}}} \right)^{\frac{D-2}{D-1}} = \sum_g C_T \left( \frac{V}{V_{\text{AdS}}} \right)^{\frac{3}{4}}$$

$\neq KS$

# Holds for all higher odd dimensions

Dimension	$\beta$ such that $\Delta\mathcal{C}_V \sim (r_+/\ell)^\beta$	$V^{(D-2)/(D-1)}$
5	4.50000	$9/2 = 4.5$
7	6.66667	$20/3 \approx 6.66667$
9	8.75000	$35/4 = 8.75$
11	10.80000	$54/5 = 10.8$
13	12.83333	$77/6 \approx 12.83333$
15	14.85714	$104/7 \approx 14.85714$
17	16.87500	$135/8 \approx 16.87500$
19	18.88889	$170/9 \approx 18.88889$
21	20.90000	$209/10 = 20.9$
23	22.90909	$252/11 \approx 22.90909$
25	24.91667	$299/12 \approx 24.91667$
27	26.92308	$350/13 \approx 26.92308$

Thermodynamic Volume

$$V = \frac{r_+^{2(N+1)} \Omega_{2N+1}}{2(N+1)} + \frac{4\pi a J}{(2N+1)(N+1)}$$

Thermodynamic Entropy

$$S = \frac{\Omega_{2N+1} h(r_+) r_+^{2N}}{4G_N}$$

# Lesson: Complexity of Formation scales with Thermodynamic Volume

Complexity = Volume

$$\Delta\mathcal{C}_V(\Sigma) = \lim_{r_{\max} \rightarrow \infty} \frac{[\mathcal{V}(\Sigma) - 2\mathcal{V}_{AdS}]}{G_N R}$$

$$\Delta\mathcal{C}_V = \Sigma_g C_T \left( \frac{V}{V_{AdS}} \right)^{\frac{D-2}{D-1}}$$

Complexity= Action

$$\Delta\mathcal{C}_A(\Sigma) = \frac{I_{WDW} - 2I_{AdS}}{\pi}$$

$$\Delta\mathcal{C}_A(\Sigma) = \Sigma_g C_T \left( \frac{V}{V_{AdS}} \right)^{\frac{D-2}{D-1}}$$

Thermodynamic Volume

$$V = \frac{r_+^{2(N+1)} \Omega_{2N+1}}{2(N+1)} + \frac{4\pi a J}{(2N+1)(N+1)}$$

Thermodynamic Entropy

$$S = \frac{\Omega_{2N+1} h(r_+) r_+^{2N}}{4G_N}$$

Also for Kerr-AdS (partial results) and BTZ!

Benamonti/Bigazzi/Billo/Faggi/Galli  
JHEP 2111 (2021) 037

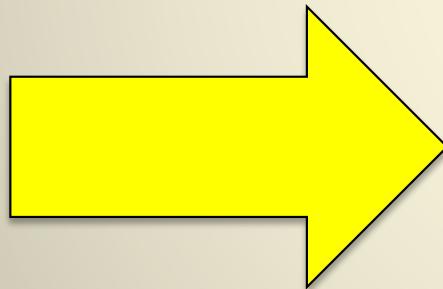
And solitons!

Andrews/Hennigar/Kunduri  
CQG 37 (2020) 204002

# Entropy Bounds Complexity

## Reverse Isoperimetric Inequality

$$\mathcal{R}^{D-2} = \left( \frac{(D-1)V}{\Omega_{D-2}} \right)^{(D-2)/(D-1)} \left( \frac{\Omega_{D-2}}{4G_N S} \right) \geq 1$$



$$\Delta \mathcal{C} \geq \beta_D S$$

Complexity scales with volume  
but is bounded from below by black hole entropy

# Central Charge Criticality

Recall Holographic Smarr Relation

$$E = TS + \tilde{\phi}\tilde{Q} + \Omega J + \tilde{\mu}C$$

Dimensional  
Scaling  
suggests

Cong/Kubiznak/Mann  
PRL 127 (2021) 091301

$$C = k \frac{l^{D-2}}{16\pi G}$$

CFT 1<sup>st</sup> Law

$$\delta E = T\delta S - pdV + \tilde{\phi}\delta\tilde{Q} + \Omega\delta J + \tilde{\mu}\delta C$$

$$\delta M = \frac{\kappa}{8\pi G} \delta A + \Omega \delta J + \phi \delta Q - \frac{V}{8\pi G} \delta \Lambda - \alpha \frac{\delta G}{G},$$

Bulk 1<sup>st</sup> Law

$$\begin{aligned} V &= \frac{8\pi G l^2}{(D-1)(D-2)} \left( M - \phi Q - (D-2)C\tilde{\mu} \right) \\ &= \frac{(D-3)}{2P} \left( \frac{D-2}{D-3} (\Omega J + TS) + \phi Q - M \right) \end{aligned}$$

$$\begin{aligned} E &= M \\ \tilde{Q} &= \frac{Ql}{\sqrt{G}} \\ \tilde{\phi} &= \frac{\phi\sqrt{G}}{l} \end{aligned}$$

Bulk Smarr  
Relation

## Bulk 1<sup>st</sup> Law

$$\delta M = \frac{\kappa}{8\pi G} \delta A + \Omega \delta J + \phi \delta Q - \frac{V}{8\pi G} \delta \Lambda - \alpha \frac{\delta G}{G},$$

Variable  $G$

$$\begin{aligned}\alpha &= \frac{1}{2} \phi Q + \tilde{\mu} C + TS \\ &= M - \Omega J - \frac{1}{2} \phi Q\end{aligned}$$

Conjugate to  $G$

Rewrite in terms  
of Pressure

$$\frac{\delta G}{G} = -\frac{2}{D} \frac{\delta C}{C} - \frac{D-2}{D} \frac{\delta P}{P}$$



$$\delta M = T \delta S + \Omega \delta J + \phi \delta Q + V_C \delta P + \mu \delta C$$

“Mixed” Bulk  
1<sup>st</sup> Law

New  
Thermodynamic  
Volume

$$V_C = \frac{2M + (D-4)\phi Q}{2DP}$$

$$\mu = \frac{2P(V_C - V)}{C(D-2)}$$

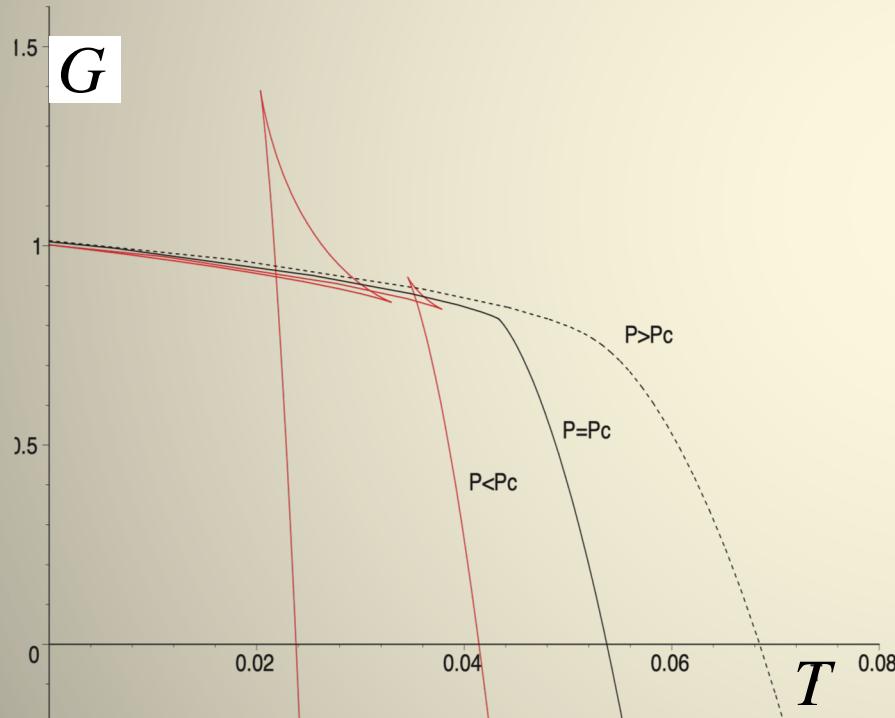
Conjugate  
Chemical  
Potential

Satisfies Reverse Isoperimetric  
Inequality for negative Chemical  
Potential

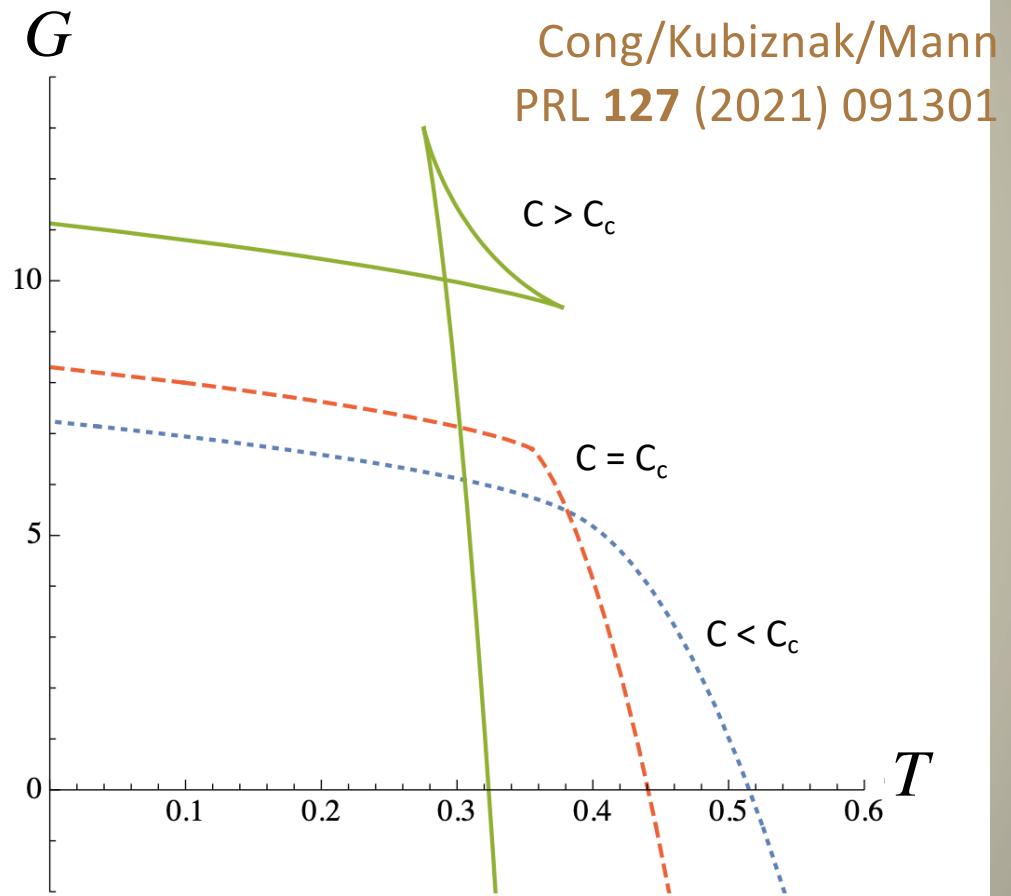
# Charged AdS Black Holes

$$G = M - TS = \frac{3GQ^2l^2 + l^2r_+^2 - r_+^4}{4Gr_+l^2} = \frac{2\pi r_+}{3}\sqrt{\frac{6PC}{k}} - \frac{2\pi Pr_+^3}{3} + \frac{3Q^2}{4r_+}$$

Kubiznak/Mann  
JHEP **1207**(2012) 033



Bulk: Critical Pressure



Boundary: Critical Central Charge

# CFT Ensembles

## No Interesting Phase Behaviour

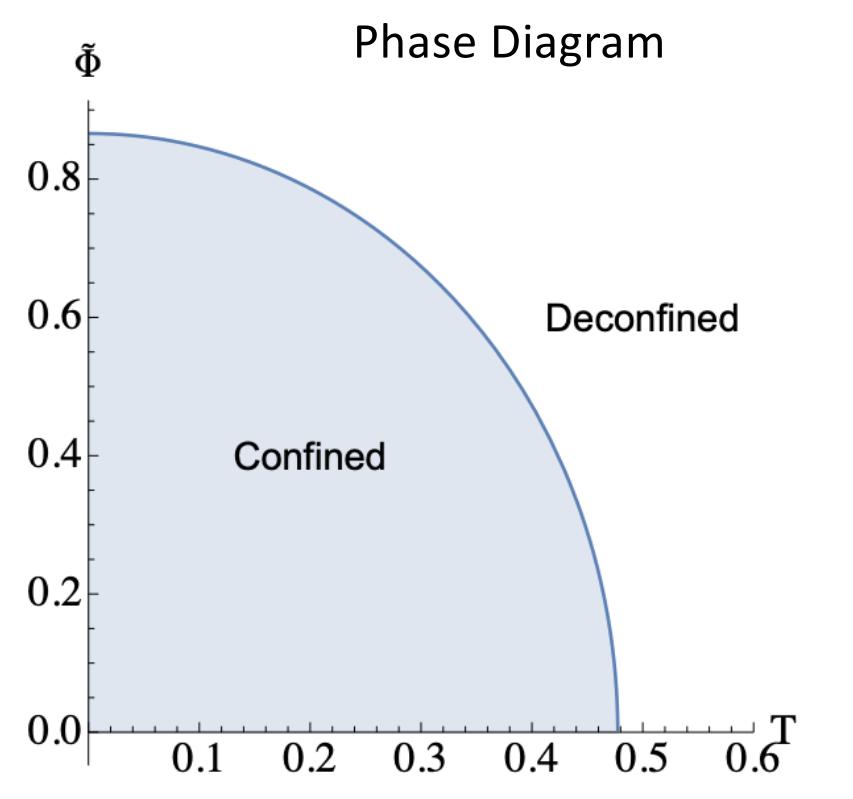
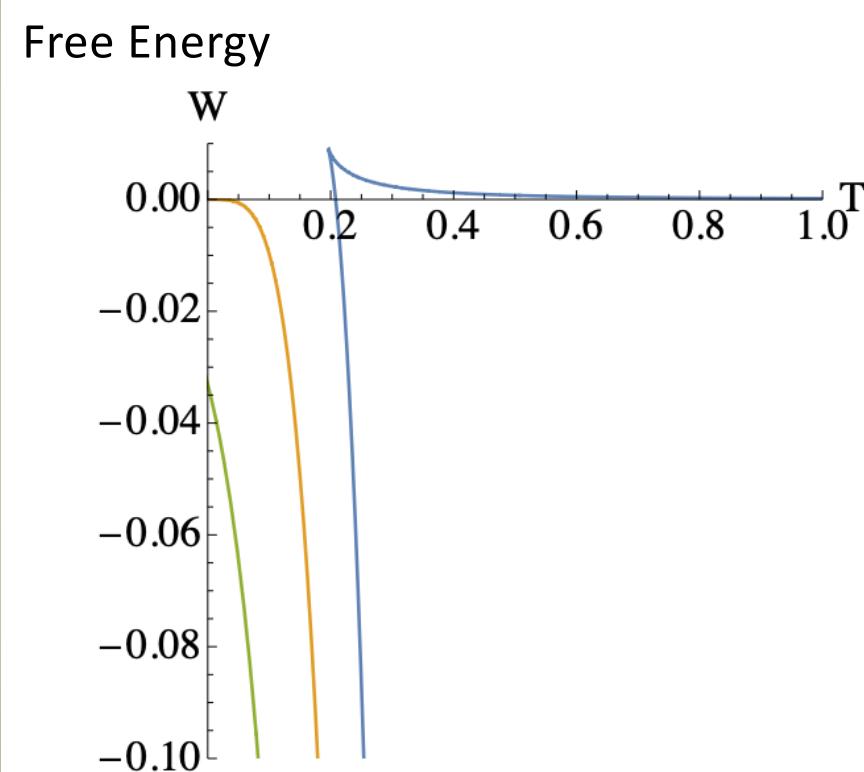
fixed $(\tilde{Q}, p, C)$ :	$F_1 \equiv E - TS + p\mathcal{V} = \tilde{\Phi}\tilde{Q} + \mu C + p\mathcal{V}$ ,
fixed $(\tilde{Q}, p, \mu)$ :	$F_2 \equiv E - TS + p\mathcal{V} - \mu C = \tilde{\Phi}\tilde{Q} + p\mathcal{V}$ ,
fixed $(\tilde{\Phi}, p, \mu)$ :	$F_3 \equiv E - TS - \tilde{\Phi}\tilde{Q} + p\mathcal{V} - \mu C = p\mathcal{V}$ ,
fixed $(\tilde{\Phi}, p, C)$ :	$F_4 \equiv E - TS - \tilde{\Phi}\tilde{Q} + p\mathcal{V} = \mu C + p\mathcal{V}$ ,
fixed $(\tilde{\Phi}, \mathcal{V}, \mu)$ :	$F_5 \equiv E - TS - \tilde{\Phi}\tilde{Q} - \mu C = 0$ .

## Interesting Phase Behaviour

fixed $(\tilde{Q}, \mathcal{V}, C)$ :	$F \equiv E - TS = \tilde{\Phi}\tilde{Q} + \mu C$ ,
fixed $(\tilde{\Phi}, \mathcal{V}, C)$ :	$W \equiv E - TS - \tilde{\Phi}\tilde{Q} = \mu C$ ,
fixed $(\tilde{Q}, \mathcal{V}, \mu)$ :	$G \equiv E - TS - \mu C = \tilde{\Phi}\tilde{Q}$

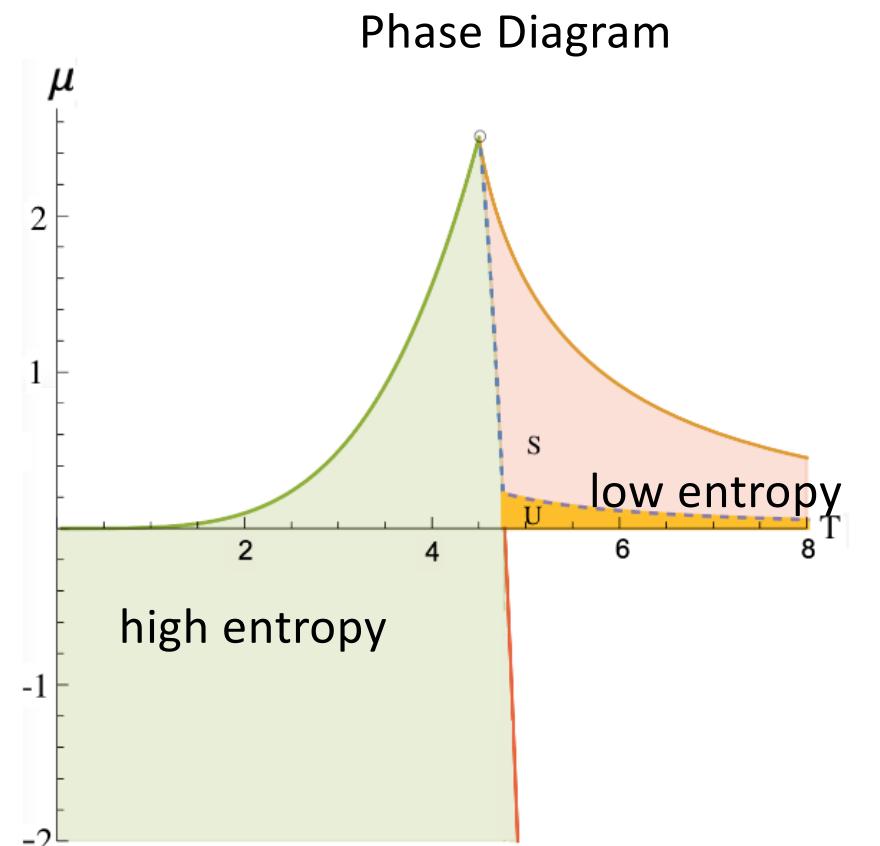
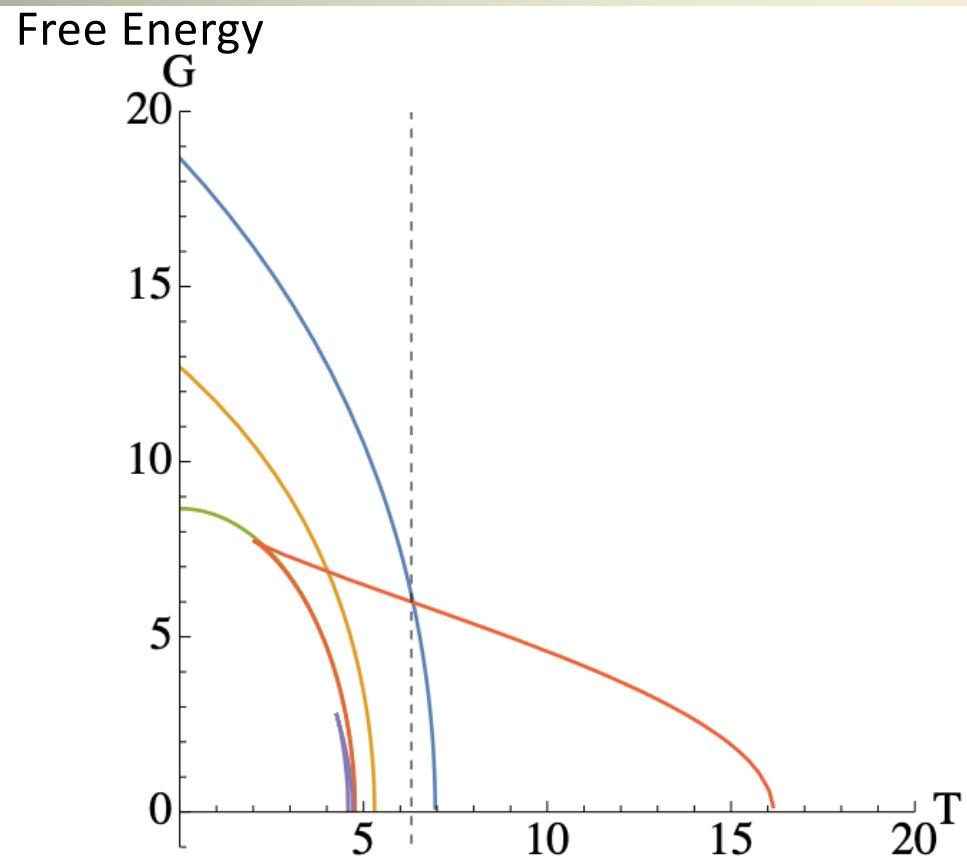
fixed  $(\tilde{\Phi}, \mathcal{V}, C)$  :

$$W \equiv E - TS - \tilde{\Phi} \tilde{Q} = \mu C,$$



fixed  $(\tilde{Q}, \mathcal{V}, \mu)$  :

$$G \equiv E - TS - \mu C = \tilde{\Phi} \tilde{Q}$$



# Holographic Dual of Black Hole Chemistry

Recall

$$C \propto \frac{L^{d-2}}{G_N}$$

Belhaj/Cong/Kubiznak/  
Mann/Visser  
PRL 130 (2023) 181401

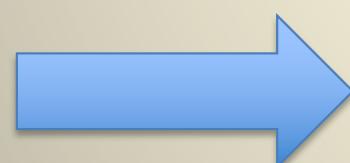
Boundary metric

$$ds^2 = \omega^2 \left( -dt^2 + L^2 d\Omega_{k,d-2}^2 \right)$$

Standard choice:

$$\omega = 1 \Rightarrow \mathcal{V} \propto L^{d-2}$$

CFT volume



$\begin{cases} \mu\delta C \text{ and } -p\delta\mathcal{V} \text{ not truly independent} \\ \text{or} \\ \text{variable } G_N \end{cases}$

Instead

$\omega$  as a thermodynamic variable

CFT volume  $\mathcal{V} \propto (\omega L)^{d-2}$

## Revised Holographic Dictionary

$$\tilde{S} = S = \frac{A}{4G_N} \quad \tilde{E} = \frac{M}{\omega} \quad \tilde{T} = \frac{T}{\omega} \quad \tilde{\Omega} = \frac{\Omega}{\omega}$$

$$\tilde{J} = J \quad \tilde{\Phi} = \frac{\Phi \sqrt{G_N}}{\omega L} \quad \tilde{Q} = \frac{QL}{\sqrt{G_N}}$$

Bulk 1<sup>st</sup> Law

$$\delta M = T\delta S + V\delta P + \Phi\delta Q + \Omega\delta J$$

Full Duality  
No Degeneracy  
 $G$  fixed

Boundary 1<sup>st</sup> Law

$$\delta \tilde{E} = \tilde{T}\delta S + \tilde{\Omega}\delta J + \tilde{\Phi}\delta \tilde{Q} + \mu\delta C - p\delta \mathcal{V}$$

$$\mu = \frac{1}{C}(\tilde{E} - \tilde{T}S - \tilde{\Omega}J - \tilde{\Phi}\tilde{Q})$$

$$p = \frac{\tilde{E}}{(d-2)\mathcal{V}}$$

Smarr  
also  
OK!

# CFT Criticality: Open Questions

- Duals of Re-entrant Phase Transitions, Triple Points, Superfluids,....
- Higher Curvature terms?
- How to incorporate  $1/N$  corrections?

# Multicritical Behaviour

- Multicritical Behaviour
    - Multiple Phases become indistinguishable at a particular set of thermodynamic parameters (P,T)
  - Do black holes exhibit such behaviour?
  - Gibbs Phase Rule
    - Constrains the number of phases in terms of the number of thermodynamic parameters

$$F = W - P + 1$$

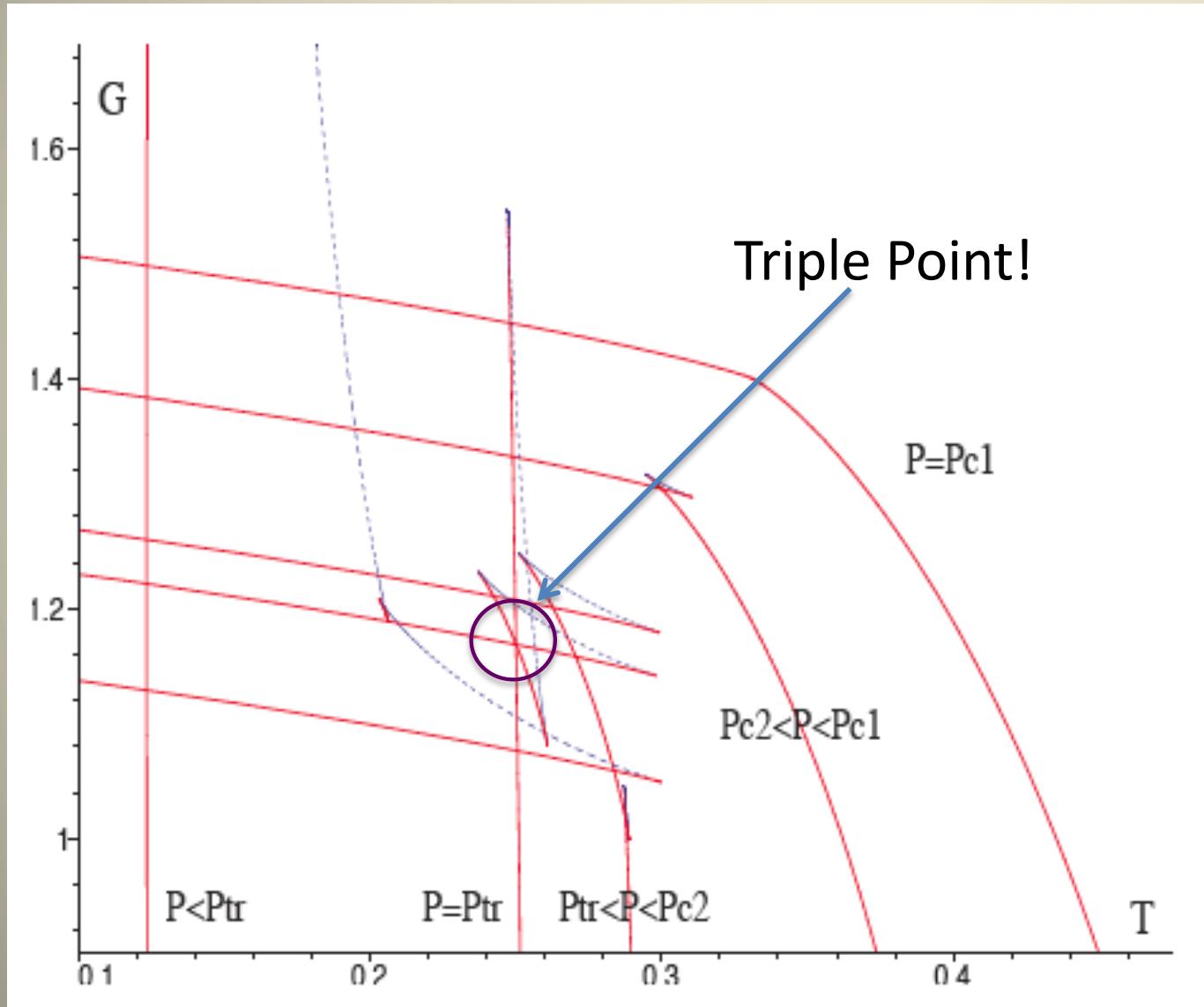
number of independent intensive parameters

number of conjugate pairs

number of coexistent phases

# Black Hole Multicriticality

Recall: Doubly rotating black hole



$$\frac{J_2}{J_1} = 0.05$$

$$P_{c1} = 0.259$$

$$P_{c2} = 0.0956$$

Altimirano/Kubiznak/  
Sherkatgnad/RBM  
CQG 31 (2014) 042001

# Multicriticality in Non-linear Electrodynamics

Tavakoli/Wu/RBM  
JHEP 2212 (2022) 117

$$S = \int d^4x \sqrt{-g} \left( R - 2\Lambda - \sum_{i=1}^N \alpha_i (F^2)^i \right)$$

$$ds^2 = -U(r)dt^2 + \frac{1}{U(r)}dr^2 + r^2d\Omega_2^2$$

$$A_\mu = [\Phi(r), 0, 0, 0]$$

$$(r(U(r) - 1))' + r^2\Lambda - r^2 \sum_{n=1}^N \left(n - \frac{1}{2}\right) \alpha_n (-2(\Phi')^2)^n = 0$$

$$\frac{1}{2}r^2 \sum_{n=1}^N n\alpha_n (-2(\Phi')^2)^n - Q(\Phi') = 0$$

$$c_1 = -2M \quad c_i = \frac{4Q}{i+2} b_{i-1}$$

$$\Phi = \sum_{i=1}^K b_i r^{-i} \quad U = 1 + \sum_{i=1}^K c_i r^{-i} + \frac{r^2}{l^2} \quad b_1 = Q \quad b_5 = \frac{4}{5} Q^3 \alpha_2$$

$$b_9 = \frac{4}{3} Q^5 (4\alpha_2^2 - \alpha_3) \quad b_{13} = \frac{32}{13} Q^7 (24\alpha_2^3 - 12\alpha_3\alpha_2 + \alpha_4)$$

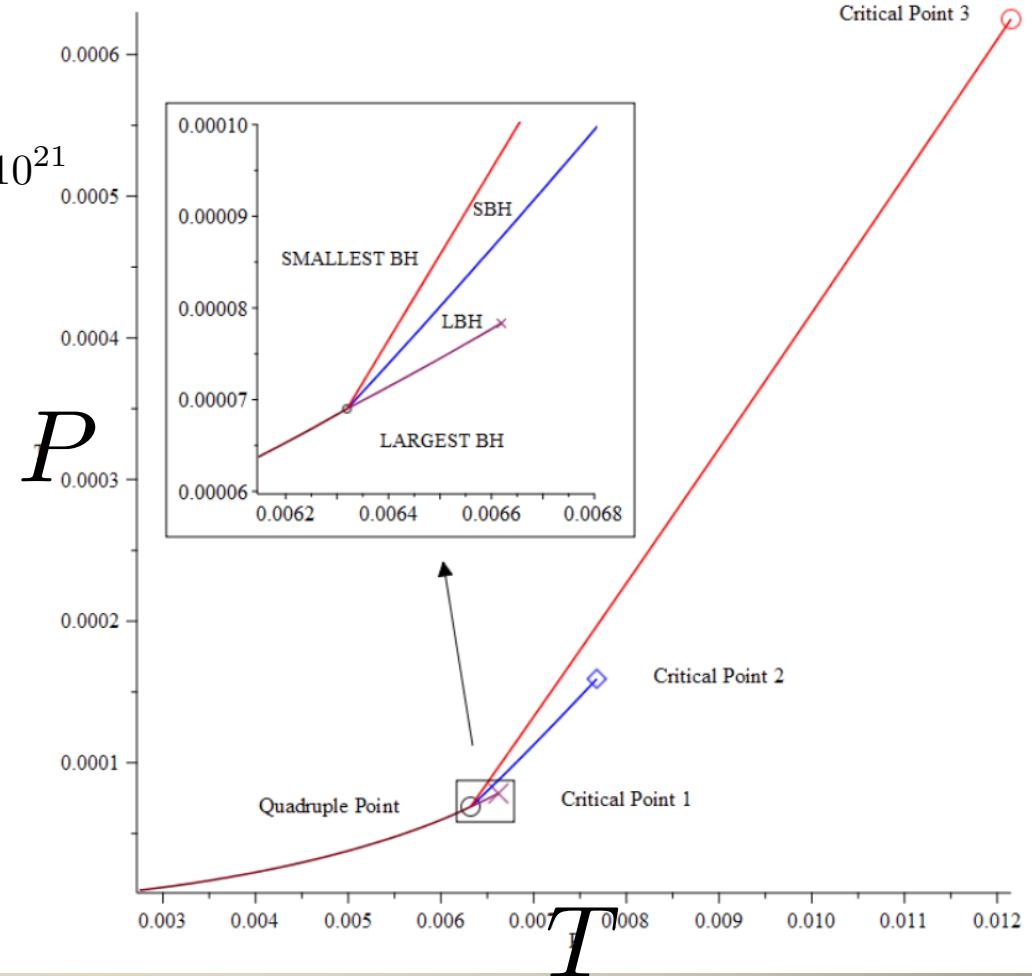
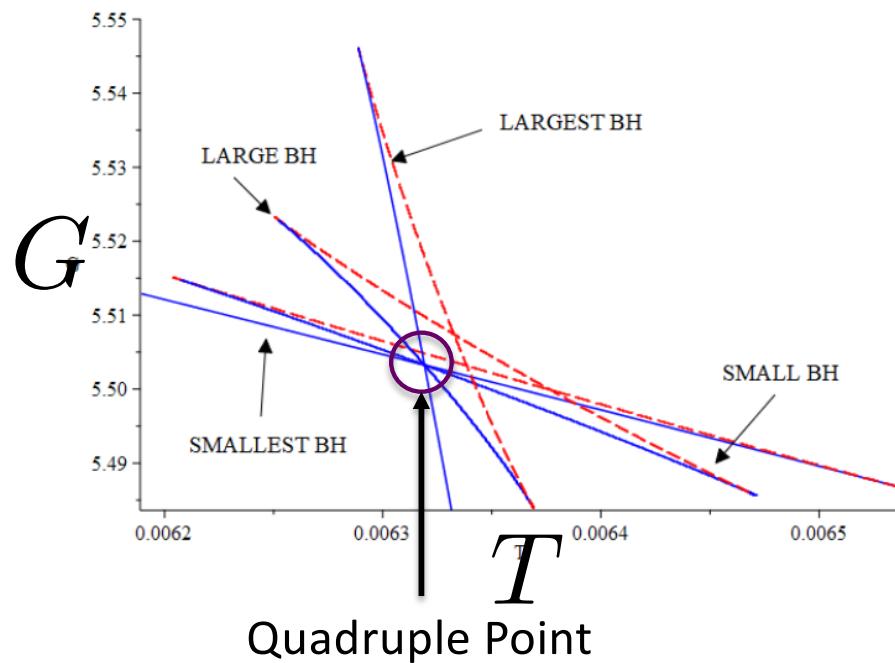
$$\begin{aligned}
P = & \frac{T}{2r_+} - \frac{1}{8\pi r_+^2} + \frac{Q^2}{8\pi r_+^4} + \frac{5b_5 Q}{16\pi r_+^8} + \frac{3b_9 Q}{8\pi r_+^{12}} \\
& + \frac{13b_{13} Q}{32\pi r_+^{16}} + \frac{17b_{17} Q}{40\pi r_+^{20}} + \frac{7b_{21} Q}{16\pi r_+^{24}} + \frac{25b_{25} Q}{56\pi r_+^{28}}
\end{aligned}$$

$$P = 6.9 \times 10^{-5} Q = 6.75112$$

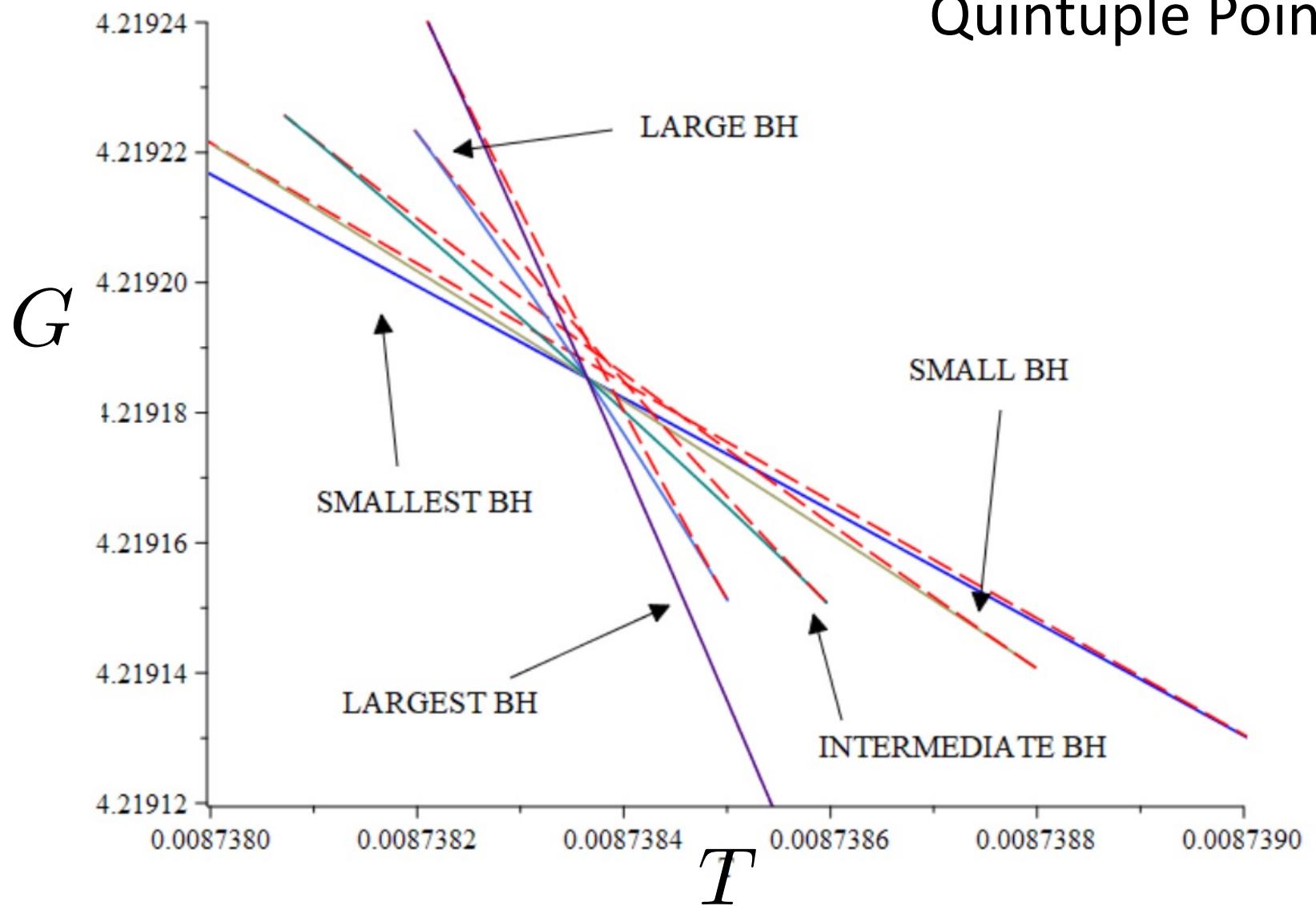
$$\alpha_2 = -20.63286635 \quad \alpha_3 = 1379.056050$$

$$\alpha_4 = -133263.0329 \quad \alpha_5 = 1.550137197 \times 10^7$$

$$\alpha_6 = -2.017480713 \times 10^9 \quad \alpha_7 = 5.046133724 \times 10^{21}$$



# Quintuple Point



# Multicriticality in Multiply Rotating Black Holes

$$ds^2 = -W \left(1 + \frac{r^2}{l^2}\right) d\tau^2 + \frac{2m}{U} \left(W d\tau - \sum_{i=1}^N \frac{a_i \mu_i^2 d\varphi_i}{\Xi_i}\right)^2$$

Kerr-AdS

$$+ \sum_{i=1}^N \frac{r^2 + a_i^2}{\Xi_i} \mu_i^2 d\varphi_i^2 + \frac{U dr^2}{F - 2m} + \sum_{i=1}^{N+\epsilon} \frac{r^2 + a_i^2}{\Xi_i} d\mu_i^2$$

$$- \frac{l^{-2}}{W(1 + r^2/l^2)} \left( \sum_{i=1}^{N+\epsilon} \frac{r^2 + a_i^2}{\Xi_i} \mu_i d\mu_i \right)^2$$

$$n = \frac{1}{2}(d - 1 - \epsilon)$$

$$W = \sum_{i=1}^{N+\epsilon} \frac{\mu_i^2}{\Xi_i} \quad U = r^\epsilon \sum_{i=1}^{N+\epsilon} \frac{\mu_i^2}{r^2 + a_i^2} \prod_j^N (r^2 + a_j^2) \quad \epsilon = 0/1 \text{ odd/even } D$$

$$\sum_{i=1}^n \mu_i^2 = 1$$

$$F = r^{\epsilon-2} \left(1 + \frac{r^2}{l^2}\right) \prod_{i=1}^N (r^2 + a_i^2) \quad \Xi_i = 1 - \frac{a_i^2}{l^2}$$

$$M = \frac{m\omega_{D-2}}{4\pi(\prod_j \Xi_j)} \left( \sum_{i=1}^N \frac{1}{\Xi_i} - \frac{1-\epsilon}{2} \right) \quad N = \frac{1}{2}(D-1-\epsilon)$$

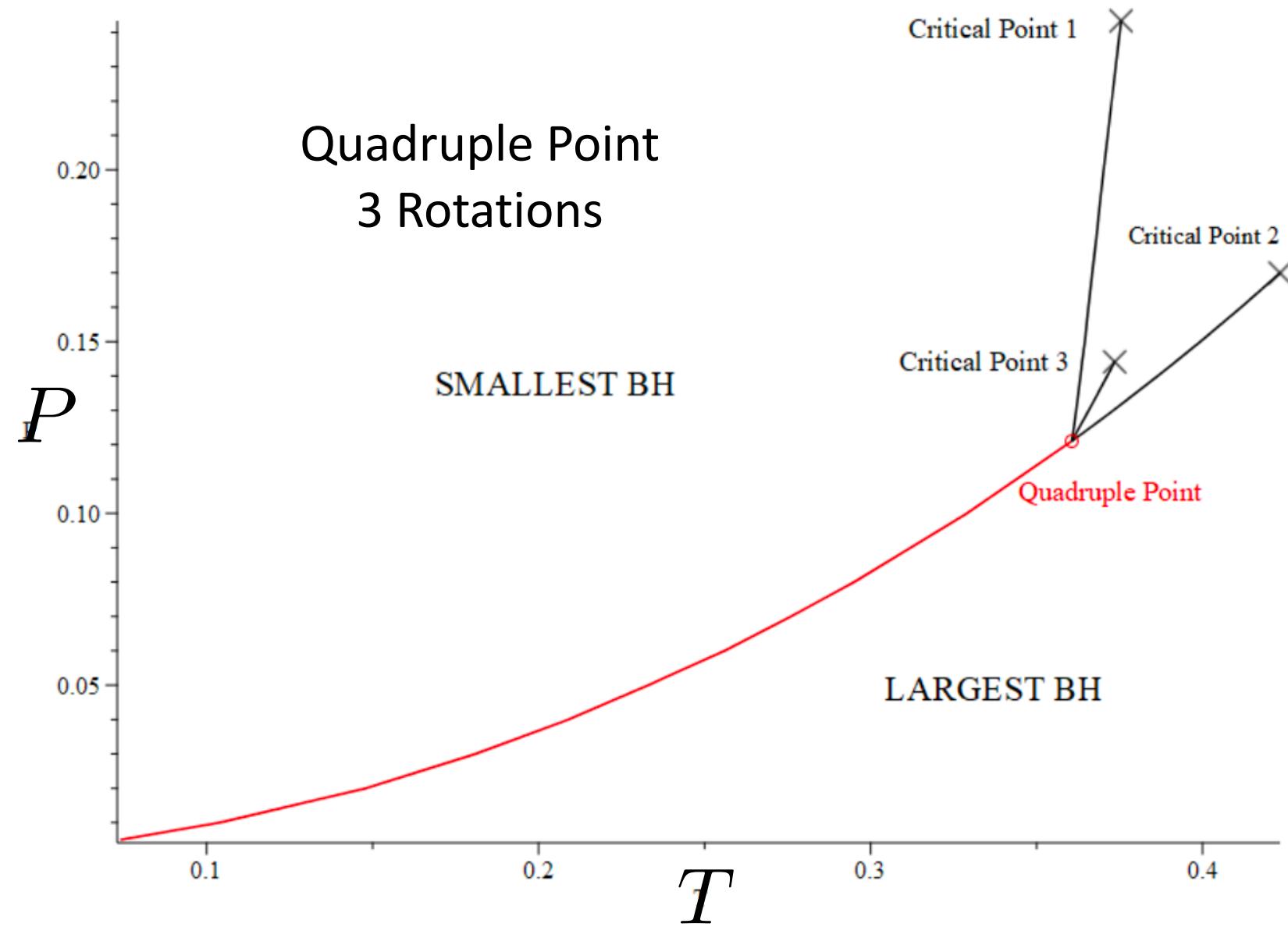
$$J_i = \frac{a_i m \omega_{D-2}}{4\pi \Xi_i (\prod_j \Xi_j)} \quad \Omega_i = \frac{a_i (1 + \frac{r_+^2}{l^2})}{r_+^2 + a_i^2}$$

$$T = \frac{1}{2\pi} \left[ r_+ \left( \frac{r_+^2}{l^2} + 1 \right) \sum_{i=1}^N \frac{1}{a_i^2 + r_+^2} - \frac{1}{r_+} \left( \frac{1}{2} - \frac{r_+^2}{2l^2} \right)^\epsilon \right]$$

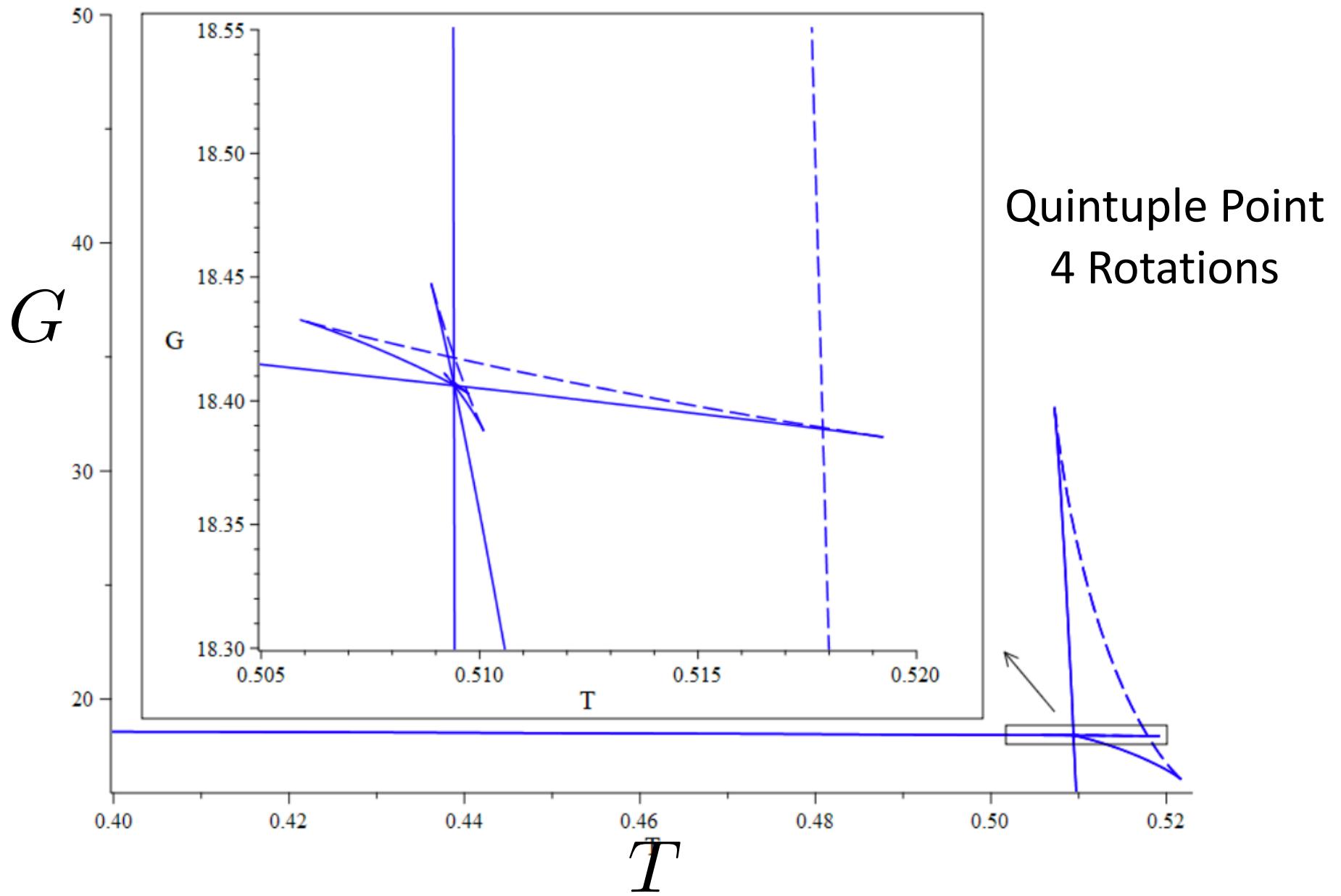
$$S = \frac{\omega_{D-2}}{4r_+^{1-\epsilon}} \prod_{i=1}^N \frac{a_i^2 + r_+^2}{\Xi_i} \quad V = \frac{r_+ A}{D-1} + \frac{8\pi}{(D-1)(D-2)} \sum_{i=1}^n a_i J_i$$

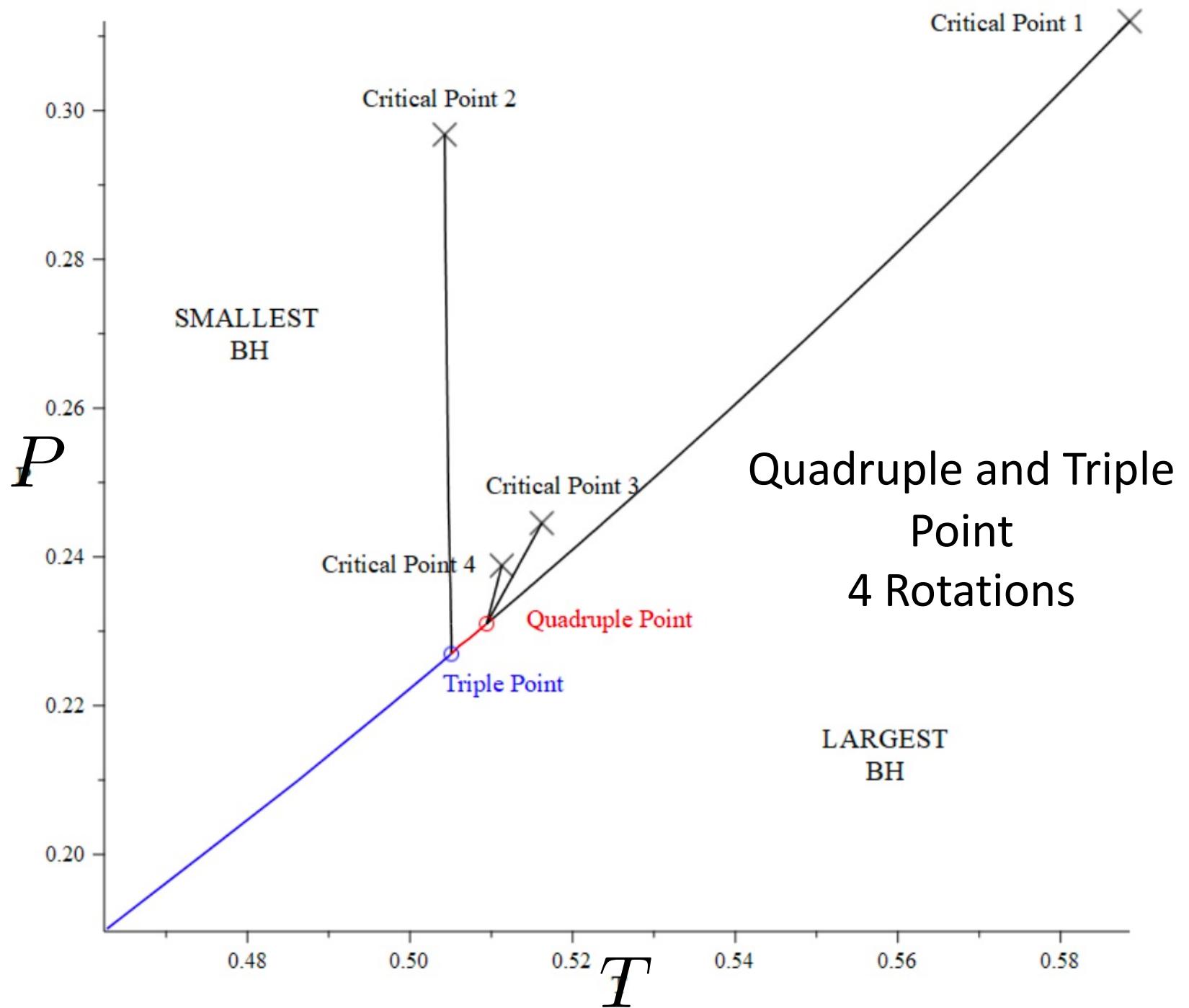
## Gibbs Phase Rule

- Two rotations: Triple Point (D > 5)
- Three rotations: Quadruple Point (D > 7)
- Four rotations: Quintuple Point (D > 9)
- Five rotations: Hextuple Point (D > 11)



$$D = 8, J_1 = 7.967, J_2 = 1.24, J_3 = 0.12798$$





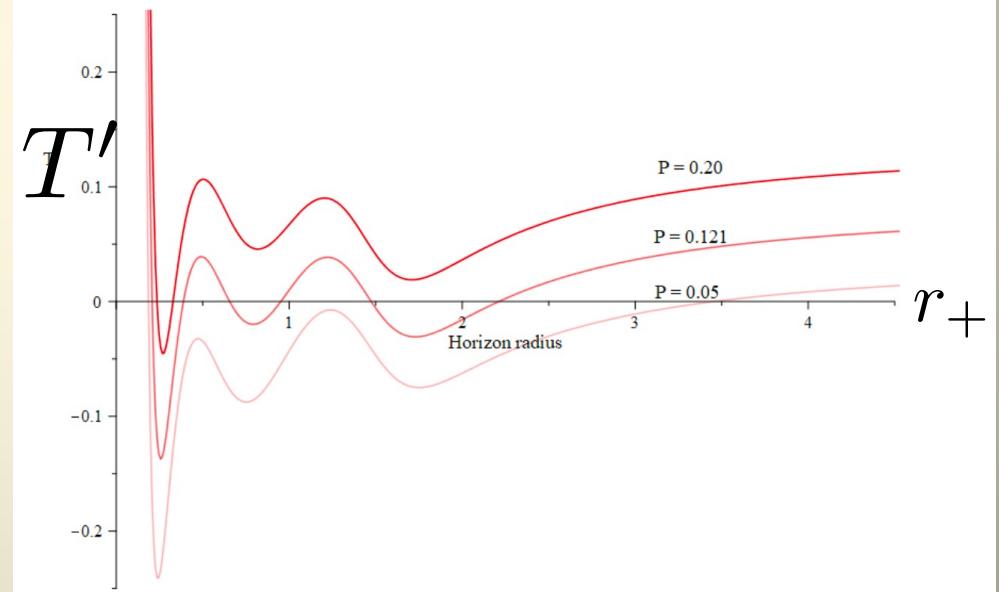
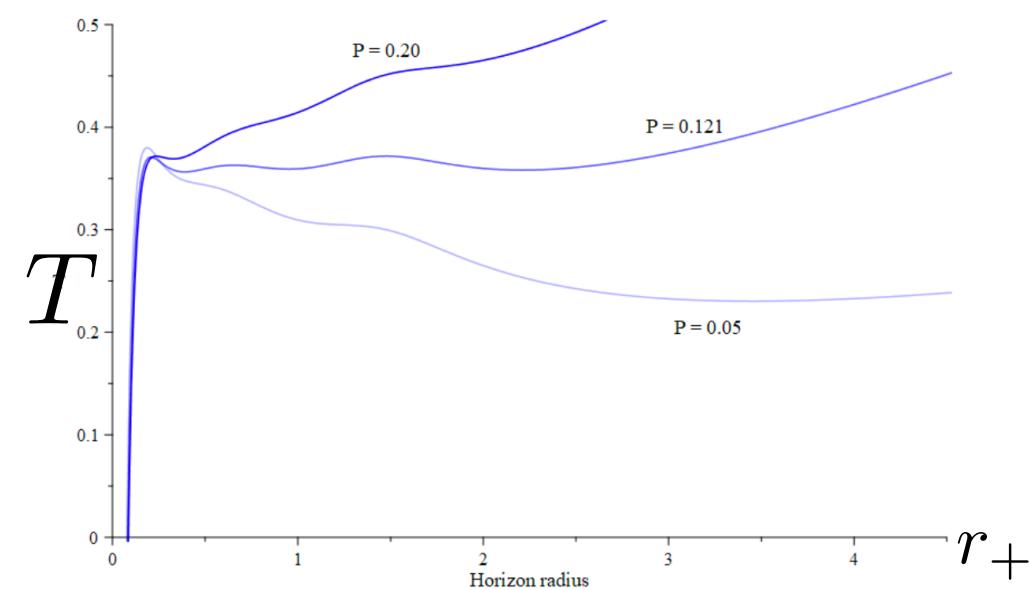
$$d = 10, J_1 = 24.48, J_2 = 4.331, J_3 = 1.1973, J_4 = 0.155$$

How?

$$dG = \sum_{i=1}^N \Omega_i dJ_i - SdT - VdP$$

→  $dG = -SdT$  for constant  $P$  and  $J_i$

→ extrema of  $G(r_+)$  and  $T(r_+)$  occur at the same  $r_+$  values

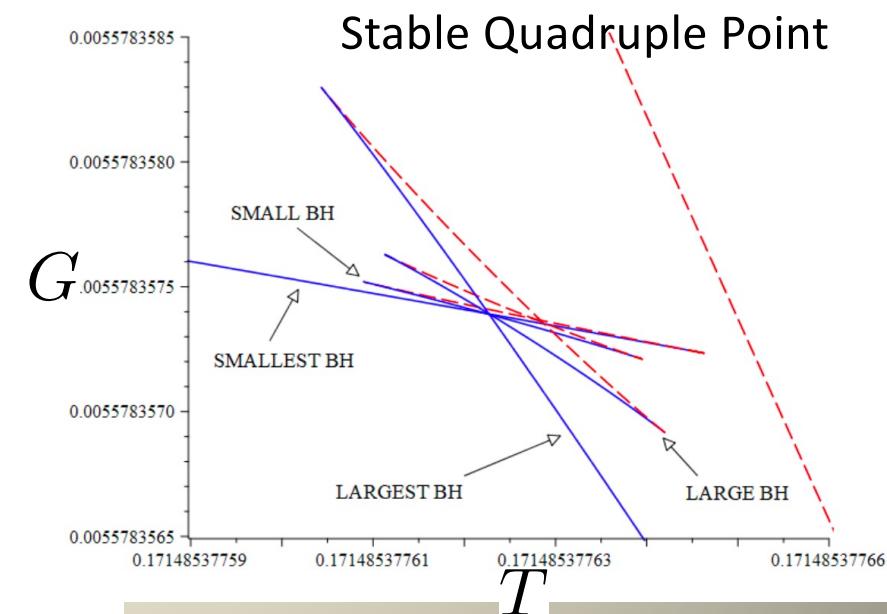
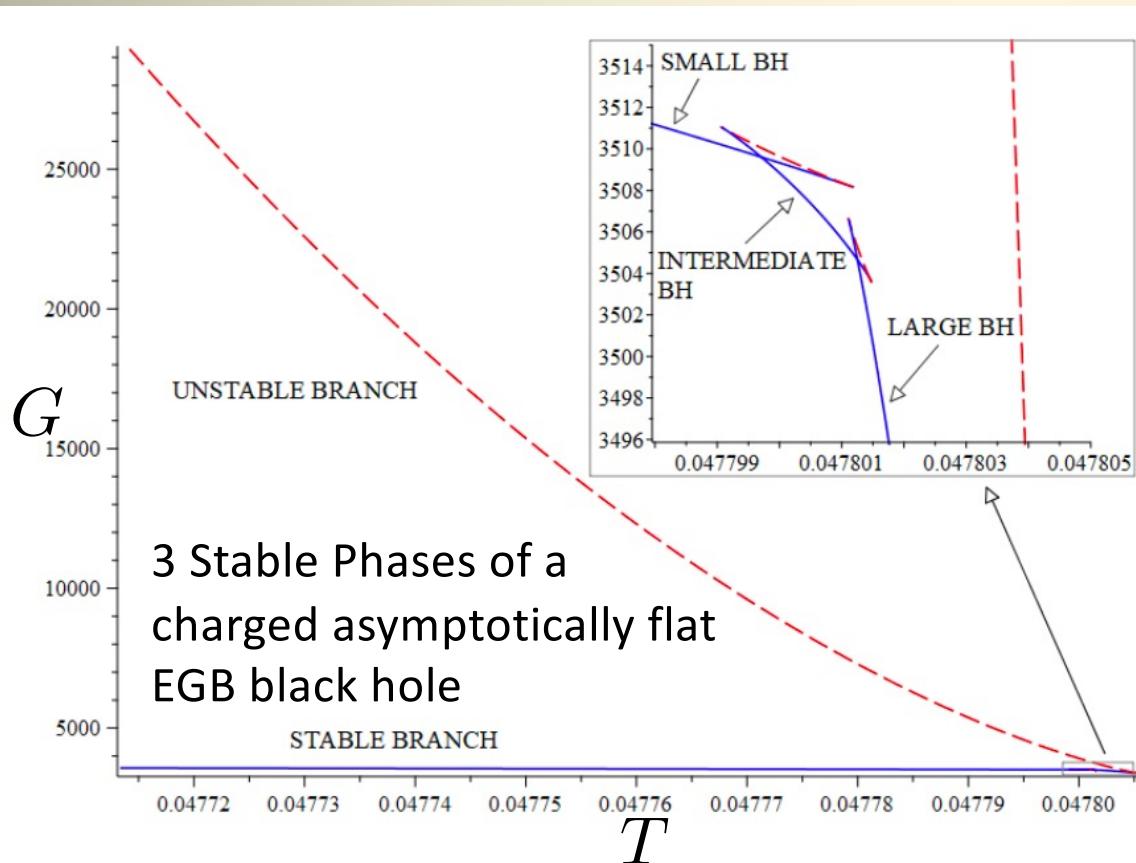


maximum number of coexistent states when  $T'(r_+)$  has a root between every local extremum

# Asymptotically Flat Black Holes

Wu/Mann  
CQG (to appear) 2212.08673

- Asymptotically Flat Black holes can exhibit multiple phases and multicritical points!



# Even more Black Hole Chemistry

- Joule-Thompson Expansion Okcu/Aydiner  
Eur.Phys.J.C **77** (2017) 1, 24
  - High-to-low pressure changes in black holes mimic those observed in real-world gases
- Thermodynamic Topological Defects Wei/Liu/Mann  
PRL **129** (2022) 191101
  - black hole solutions as defects in thermodynamic parameter space
  - winding numbers indicate stability
- Dynamics of Phase Transitions Li/Wang PRD **102** (2020) 024085  
Li/Zhang/Wang JHEP **2010** (2020) 09
  - Can use Smoluchowski Eqn to probe dynamics of phase changes
  - Works at multicritical points

# Take Home Lesson

- Thermodynamics is Thermodynamics
  - Black Holes exhibit pretty much every behaviour seen in chemistry under the right circumstances
  - Lesson for Quantum Gravity?
- AdS is not a restriction
  - Chemical Phenomena seen in asymptotically flat and asymptotically de Sitter settings
- Holographic understanding emerging
- Still a lot to learn from Black Hole Chemistry

Kubiznak/Mann/Teo,  
*Black hole chemistry: Thermodynamics with Lambda*, CQG 34 (2017) 063001